

Active phase separation in mixtures of chemically interacting particles

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MAX-PLANCK-GESELLSCHAFT



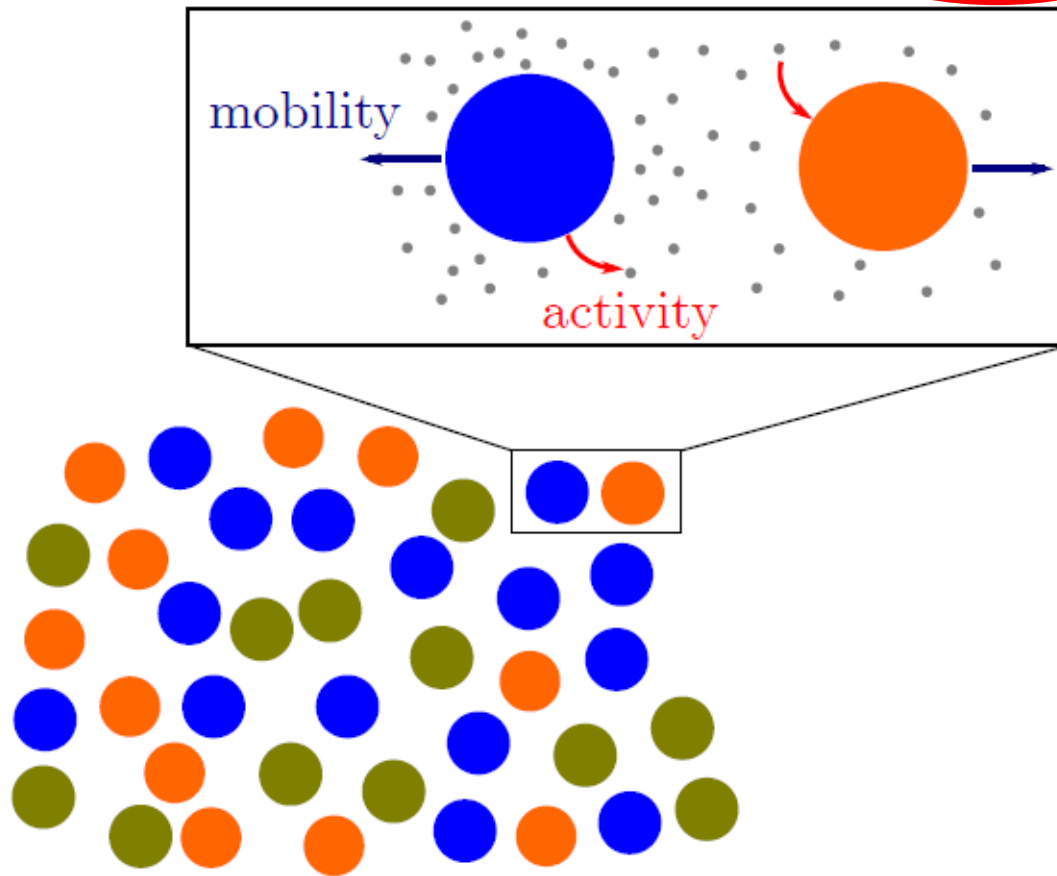
PennState



Chemically interacting particles:

microorganisms, catalytic colloids, enzymes..

???

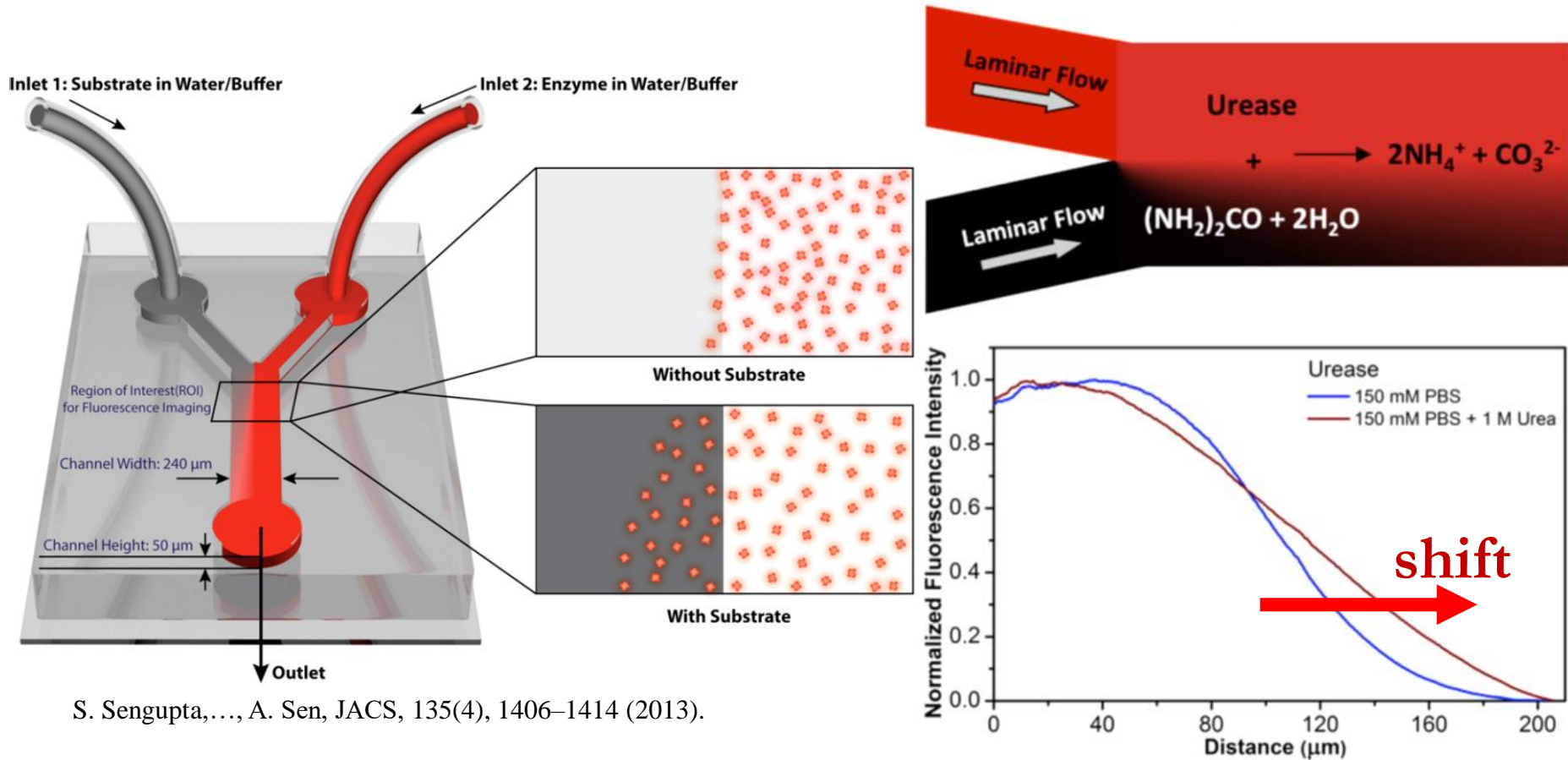


► will such a mixture remain homogeneous?

Detour:

**Chemotaxis of catalytic enzymes
in the presence of their substrate**

Experimental Observations of Enzyme Chemotaxis



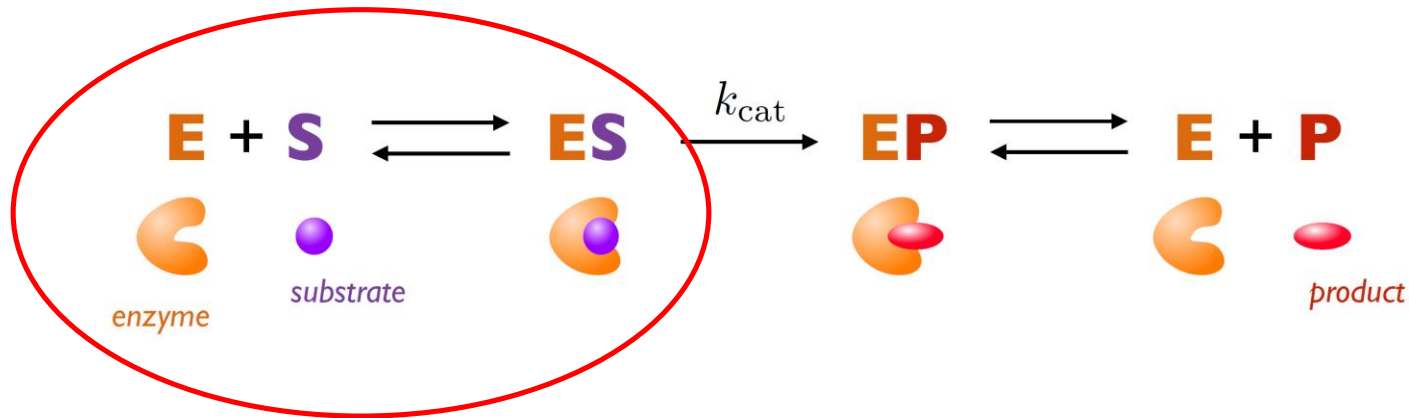
Typically **towards** the substrate, occasionally **away** from it

Experimental Observations of Enzyme Chemotaxis

- Yu et al. JACS, 131(16), 5722–5723 (2009)
RNA polymerase: Chemotaxis
- Sengupta et al. JACS 135(4), 1406–1414 (2013)
Catalase, urease: Chemotaxis
- Sengupta et al. ACS Nano, 8(3), 2410–2418 (2014)
DNA polymerase: Chemotaxis
- Zhao et al. Nat. Chem, 10, 311–317 (2018)
Hexokinase, aldolase: Chemotaxis
- Jee et al. PNAS, 139(44), 15588–15591 (2018)
Urease, acetylcholinesterase: Antichemotaxis (!)

Conflicting observations for urease!

What is the mechanism of enzyme chemotaxis?

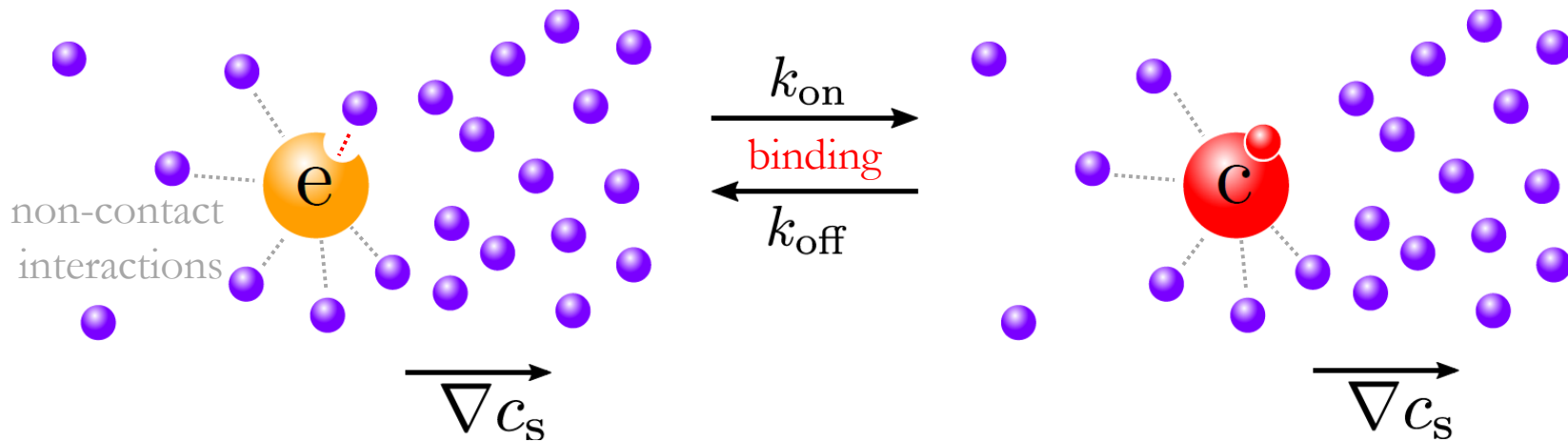


- **Active** mechanisms (self-phoresis, stochastic swimming...) depending on the catalytic step seem to be **too weak to explain observations**

[Golestanian, R. (2015). *Phys. Rev. Lett.*, 115(10), 108102.]

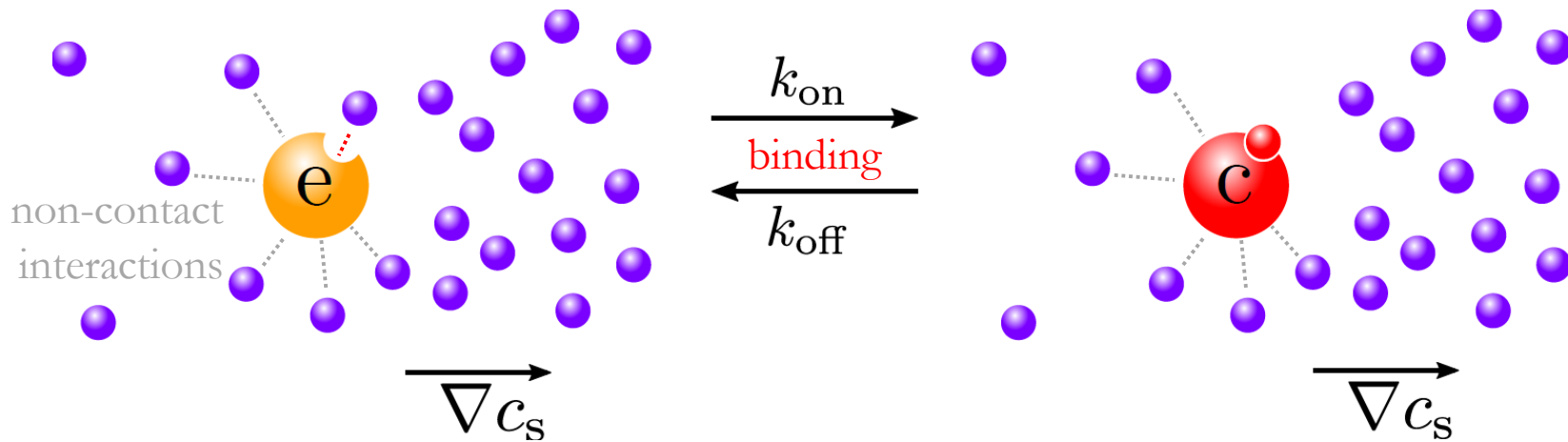
- Could **passive** mechanisms (just binding-unbinding) explain the observations?

A Microscopic Model of Enzyme Chemotaxis



- Interactions between enzyme and substrate:
 - Binding to form a complex
 - Non-contact (electrostatic, van der Waals, etc.)
 - Hydrodynamic

A Microscopic Model of Enzyme Chemotaxis



- Starting from the full Fokker-Planck equation for the $(N+1)$ -particle probability distribution, integrating out the substrate we obtain:

$$\partial_t c_e(\mathbf{R}; t) = \nabla_{\mathbf{R}} \cdot [D_e \nabla_{\mathbf{R}} c_e - \mathbf{v}_e(\mathbf{R}) c_e] - k_{\text{on}} c_e c_s + k_{\text{off}} c_c$$

$$\partial_t c_c(\mathbf{R}; t) = \nabla_{\mathbf{R}} \cdot [D_c \nabla_{\mathbf{R}} c_c - \mathbf{v}_c(\mathbf{R}) c_c] + k_{\text{on}} c_e c_s - k_{\text{off}} c_c$$

↑
↑
↑
↑

diffusion
phoresis
binding
unbinding

A Microscopic Model of Enzyme Chemotaxis

$$\partial_t c_e(\mathbf{R}; t) = \nabla_{\mathbf{R}} \cdot [D_e \nabla_{\mathbf{R}} c_e - \mathbf{v}_e(\mathbf{R}) c_e] - k_{\text{on}} c_e c_s + k_{\text{off}} c_c$$

$$\partial_t c_c(\mathbf{R}; t) = \nabla_{\mathbf{R}} \cdot [D_c \nabla_{\mathbf{R}} c_c - \mathbf{v}_c(\mathbf{R}) c_c] + k_{\text{on}} c_e c_s - k_{\text{off}} c_c$$

Diffusion coefficients of free enzyme and complex different in general $D_e \neq D_c$ (**conformational changes!**)

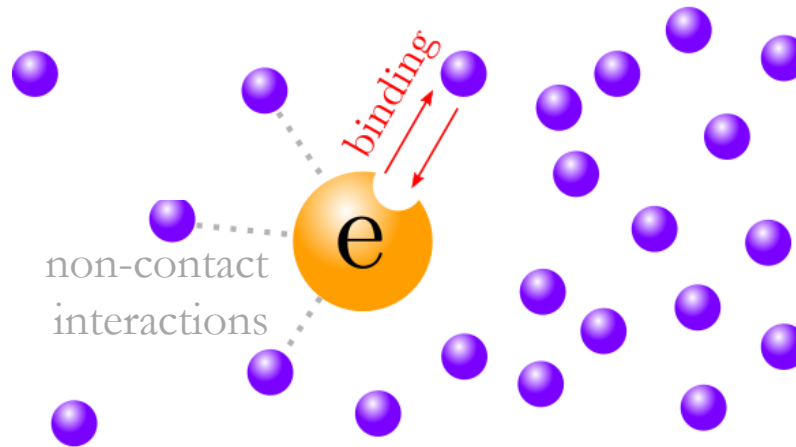
Phoresis due to **non-contact** interactions with substrate:

$$\mathbf{v}_e(\mathbf{R}) \approx \frac{k_B T}{\eta} \left[\int_0^\infty dh h \left(e^{-\phi^{\text{es}}(h)/k_B T} - 1 \right) \right] \nabla_{\mathbf{R}} c_s \equiv \frac{k_B T}{\eta} \lambda_e^2 \nabla_{\mathbf{R}} c_s$$

$$\mathbf{v}_c(\mathbf{R}) \approx \frac{k_B T}{\eta} \left[\int_0^\infty dh h \left(e^{-\phi^{\text{cs}}(h)/k_B T} - 1 \right) \right] \nabla_{\mathbf{R}} c_s \equiv \frac{k_B T}{\eta} \lambda_c^2 \nabla_{\mathbf{R}} c_s$$

Derjaguin lengths of free enzyme and complex different in general $\lambda_e \neq \lambda_c$

A Microscopic Model of Enzyme Chemotaxis



Experiments with fluorescently-labelled enzymes measure the **total enzyme concentration**

$$c_e^{\text{tot}}(\mathbf{R}; t) = c_e(\mathbf{R}; t) + c_c(\mathbf{R}; t)$$

Binding-unbinding much faster than diffusion;
we have **instantaneous local binding equilibrium**

$$k_{\text{on}}c_e(\mathbf{R}; t)c_s(\mathbf{R}; t) \approx k_{\text{off}}c_c(\mathbf{R}; t)$$

Diffusion and Chemotaxis of an Enzyme

We finally obtain:

$$\partial_t c_e^{\text{tot}}(\mathbf{R}; t) = \nabla_{\mathbf{R}} \cdot \left\{ D(\mathbf{R}) \cdot \nabla_{\mathbf{R}} c_e^{\text{tot}} - [\mathbf{V}_{\text{ph}}(\mathbf{R}) + \mathbf{V}_{\text{bi}}(\mathbf{R})] c_e^{\text{tot}} \right\}$$

Substrate-dependent **diffusion coefficient**

$$D(\mathbf{R}) = D_e + (D_c - D_e) \frac{c_s(\mathbf{R})}{K + c_s(\mathbf{R})} \quad \text{With Michaelis-Menten style dependence...}$$

Substrate-dependent **phoretic velocity**

$$\mathbf{V}_{\text{ph}}(\mathbf{R}) = \mathbf{v}_e(\mathbf{R}) + [\mathbf{v}_c(\mathbf{R}) - \mathbf{v}_e(\mathbf{R})] \frac{c_s(\mathbf{R})}{K + c_s(\mathbf{R})}$$

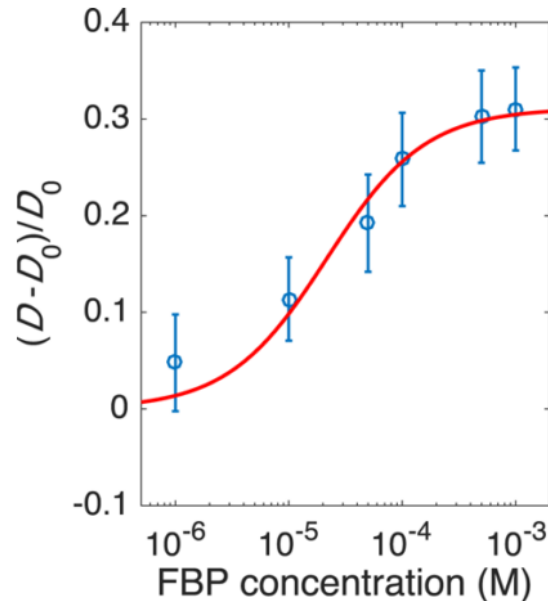
New **binding-induced contribution to chemotaxis**

$$\mathbf{V}_{\text{bi}}(\mathbf{R}) = -(D_c - D_e) \nabla_{\mathbf{R}} \left(\frac{c_s(\mathbf{R})}{K + c_s(\mathbf{R})} \right)$$

Aside: Enhanced Diffusion

It has been observed experimentally that most enzymes exhibit **enhanced diffusion** in the presence of their substrate

Example: Aldolase



$$D(\mathbf{R}) = D_e + (D_c - D_e) \frac{c_s(\mathbf{R})}{K + c_s(\mathbf{R})}$$

with $D_c > D_e$

associated to binding-induced changes in the **average shape and shape fluctuation spectrum of enzymes**

P. Illien et al., Nano Lett. 17(7), 4415–4420. (2017)

P. Illien, T. Adeleke-Larodo, and R. Golestanian, EPL 119, 40002. (2017)

J.A.-C., T. Adeleke-Larodo, P. Illien, & R. Golestanian, Acc. Chem. Res. 51, 2365. (2018)

Relating Enhanced Diffusion and Chemotaxis

Binding-induced contribution to chemotaxis can be rewritten as

$$\mathbf{V}_{\text{bi}}(\mathbf{R}) = -(D_c - D_e) \nabla_{\mathbf{R}} \left(\frac{c_s(\mathbf{R})}{K + c_s(\mathbf{R})} \right) = -\nabla_{\mathbf{R}} D(\mathbf{R})$$

If **enhanced** diffusion $D_c > D_e \Rightarrow \mathbf{V}_{\text{bi}}$ **points away from substrate**

If **hindered** diffusion $D_c < D_e \Rightarrow \mathbf{V}_{\text{bi}}$ **points towards substrate**

However, there is also the phoretic contribution...

Competition between Enhanced Diffusion and Phoresis

Two distinct chemotactic mechanisms:

1. Phoresis due to non-contact interactions

$$\mathbf{V}_{\text{ph}} \approx \frac{k_{\text{B}}T}{\eta} \lambda_{\text{e}}^2 \nabla_{\mathbf{R}} c_{\text{s}} \quad (\lambda_{\text{c}} \approx \lambda_{\text{e}})$$

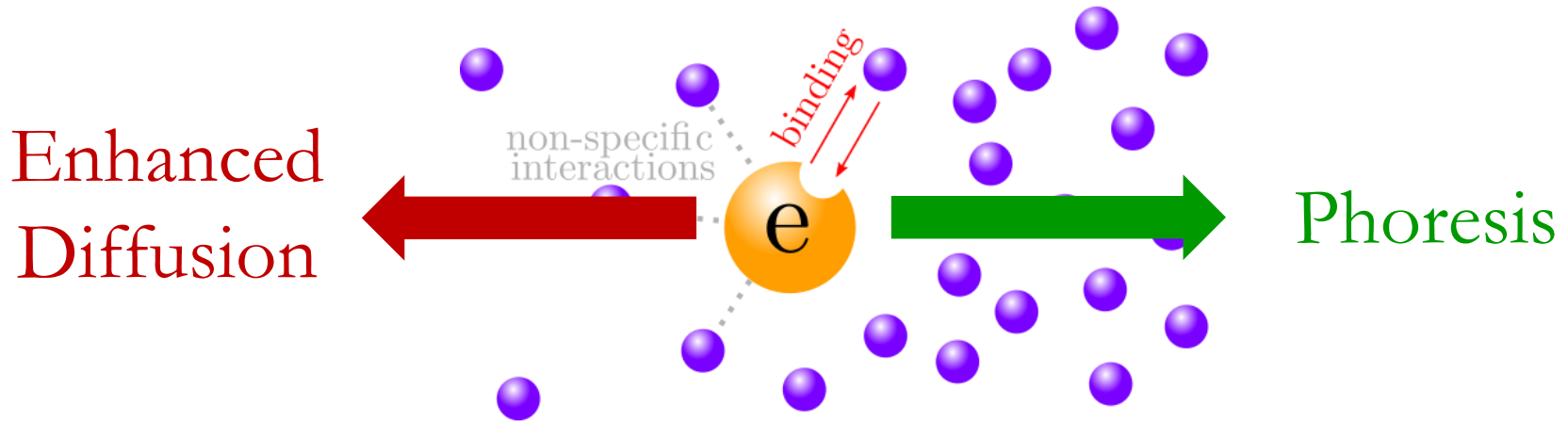
Typically **towards substrate** because interactions are predominantly attractive

2. Binding-induced changes in diffusion

$$\mathbf{V}_{\text{bi}} = -(D_{\text{c}} - D_{\text{e}}) \nabla_{\mathbf{R}} \left(\frac{c_{\text{s}}(\mathbf{R})}{K + c_{\text{s}}(\mathbf{R})} \right)$$

Typically **away from substrate** because enzymes display enhanced diffusion

Competition between Enhanced Diffusion and Phoresis



$$\mathbf{V}_{\text{bi}} = -(D_c - D_e) \frac{K}{(K + c_s)^2} \nabla_{\mathbf{R}} c_s$$

Magnitude decreases with c_s

$$\mathbf{V}_{\text{ph}} \approx \frac{k_B T}{\eta} \lambda_e^2 \nabla_{\mathbf{R}} c_s$$

~Independent of c_s

Therefore there is a **critical substrate concentration c_s^***

$c_s < c_s^* \Rightarrow$ Enhanced diffusion dominates

$c_s > c_s^* \Rightarrow$ Phoresis dominates

Making Sense of Urease Chemotaxis

Sengupta et al., JACS (2013):

Chemotaxis

$$c_s = 1 \text{ M}$$

Dominated by **phoresis**?

Jee et al., PNAS (2018):

Antichemotaxis

$$c_s = 1 \text{ mM}$$

Dominated by **enhanced diff.**?

Critical concentration:

$$c_s^* = \sqrt{\frac{K|\alpha|}{6\pi R_e |\lambda_e^2|}} - K$$

$$[\alpha \equiv (D_c - D_e)/D_e]$$

All parameters are known from experiment except for the **Derjaguin length λ_e**

$$(K = 3 \text{ mM}, \alpha = 0.3, R_e = 7 \text{ nm})$$

The two experiments will be dominated by the two different mechanisms if:

$$1 \text{ M} > c_s^* > 1 \text{ mM} \Leftrightarrow \mathbf{0.04 \text{ \AA} < \lambda_e < 8 \text{ \AA}}$$

...which is within the expected range of values

Beyond Enzymes: Designing Directed Nanovehicles

- Exploiting the competition between mechanisms:

Enhanced Diffusion

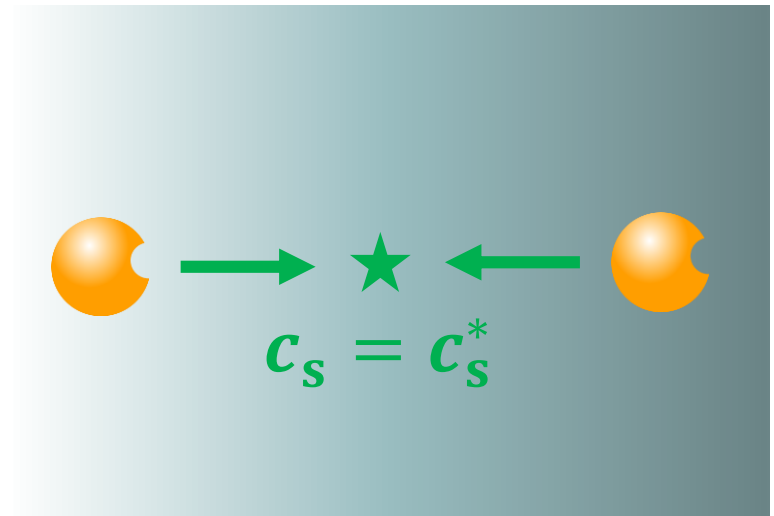
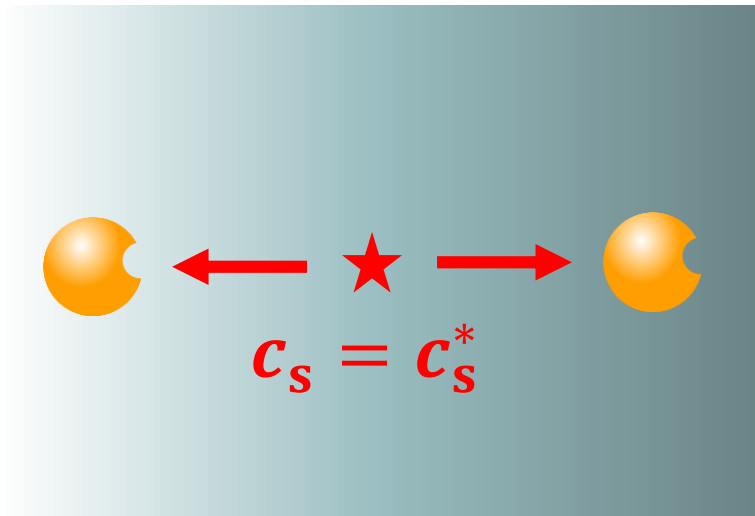
+

Attractive Phoresis:
'repelled' from regions
with $c_s = c_s^*$

Inhibited Diffusion

+

Repulsive Phoresis:
'attracted' to regions
with $c_s = c_s^*$



Summary I:

- **Enzymes undergo chemotaxis** in the presence of gradients of their substrate
- **At least two different mechanisms** contribute to chemotaxis
- Competition between phoresis and binding-induced changes in diffusion can **explain conflicting observations** for urease
- Competition between mechanisms can be exploited to engineer nanovehicles with **finely-tuned directed motion**

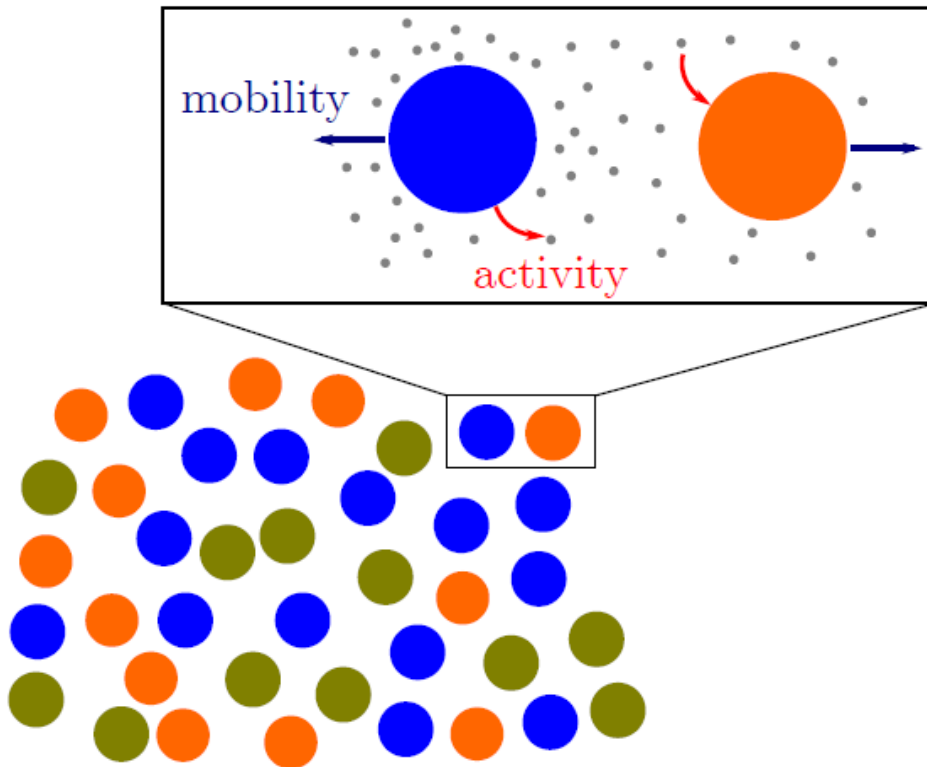
[Agudo-Canalejo, J., Illien, P., & Golestanian, R. (2018). Phoresis and Enhanced Diffusion Compete in Enzyme Chemotaxis. *Nano Letters*, 18(4), 2711–2717.]

Back on track:

**Active phase separation in mixtures
of chemically interacting particles**

Chemically-interacting particles:

microorganisms, catalytic colloids, enzymes...

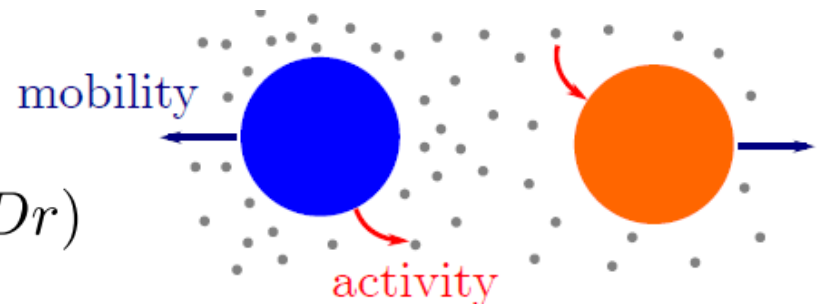


- $i = 1, \dots, M$
different particle species
- α_i chemical activity
 - positive if emits solute
 - negative if consumes solute
- μ_i chemotactic mobility
 - positive if repelled from solute
 - negative if attracted to solute

► will such a mixture remain homogeneous?

Chemical interactions are generally **non-reciprocal**

- Concentration of chemical around a single particle: $c(r) = c_0 + \alpha/(4\pi Dr)$



- Velocity response to gradient: $\mathbf{V} = -\mu \nabla c$
- Velocity of particle j in the presence of particle i : $V_{ij} \propto \alpha_i \mu_j \mathbf{r}_{ij} / |\mathbf{r}_{ij}|^3$
Velocity of i in the presence of j : $V_{ji} \propto -\alpha_j \mu_i \mathbf{r}_{ij} / |\mathbf{r}_{ij}|^3$
- Generally $V_{ij} \neq -V_{ji}$ because $\alpha_i \mu_j \neq \alpha_j \mu_i$
 \Rightarrow **non-reciprocal interactions!**
- **Action \neq reaction, signature of non-equilibrium activity**

Continuum theory for chemically-interacting mixtures

- Concentration field of particles of species $i = 1, \dots, M$:

$$\partial_t \rho_i(\mathbf{r}, t) - \nabla \cdot [D_p \nabla \rho_i + (\mu_i \nabla c) \rho_i] = 0$$

- Concentration field of chemical:

$$\partial_t c(\mathbf{r}, t) - D \nabla^2 c = \sum_i \alpha_i \rho_i$$

- Deviation from homogeneous solution:

$$\rho_i(\mathbf{r}, t) = \rho_{0i} + \delta \rho_i(\mathbf{r}, t)$$

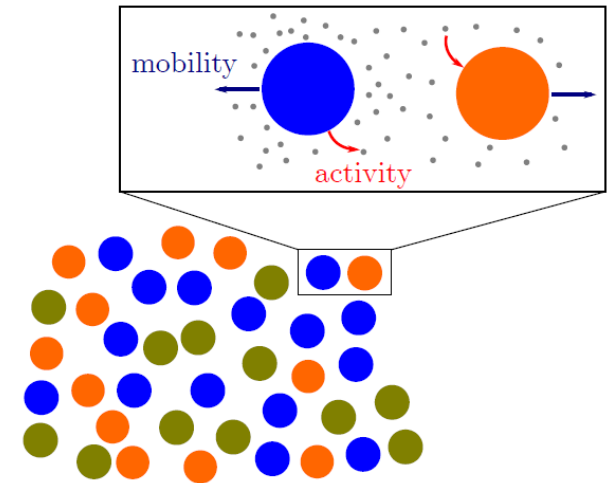
$$c(\mathbf{r}, t) = c_0 + At + \delta c(\mathbf{r}, t) \quad \text{with} \quad A \equiv \sum_i \alpha_i \rho_{0i}$$

- Equation for chemical can be rewritten as:

$$\partial_t \delta c(\mathbf{r}, t) - D \nabla^2 \delta c = \sum_i \alpha_i \delta \rho_i$$

- Fast-diffusing chemical \rightarrow instantaneous steady state

$$D \nabla^2 \delta c = - \sum_i \alpha_i \delta \rho_i$$



Continuum theory for chemically-interacting mixtures

- Concentration field of particles of species i to **linear order**:

$$\partial_t \delta \rho_i(\mathbf{r}, t) = D_p \nabla^2 \delta \rho_i - \frac{\mu_i \rho_{0i}}{D} \sum_j \alpha_j \delta \rho_j$$

- M coupled equations $\rightarrow M \times M$ eigenvalue problem in Fourier space
- We find $M - 1$ identical eigenvalues which are always negative
 $\lambda_- = -D_p q^2 \rightarrow$ corresponding modes are always stable
- ... and one eigenvalue which can become positive

$$\lambda_+ = -D_p q^2 - \sum_i \mu_i \alpha_i \rho_{0i} / D \rightarrow \text{corresponding mode can become unstable!}$$

Continuum theory: **macroscopic instability**

$$\lambda_+ = -D_p q^2 - \sum_i \mu_i \alpha_i \rho_{0i} / D$$

$$\sum_i \mu_i \alpha_i \rho_{0i} < 0$$

- **Instability** occurs whenever

- The modes with $q^2 < -\frac{\sum_i \mu_i \alpha_i \rho_{0i}}{D D_p}$ are unstable

- **Macroscopic phase separation** with $q = 0$ as the first and most unstable mode

- The corresponding eigenvectors can be calculated as

$$(\delta \rho_1, \delta \rho_2, \dots, \delta \rho_M) = \left(1, \frac{\mu_2 \rho_{02}}{\mu_1 \rho_{01}}, \dots, \frac{\mu_M \rho_{0M}}{\mu_1 \rho_{01}} \right) \delta \rho_1$$

and tell us about the **stoichiometry** at the onset of the instability

Continuum theory: **single species**

- Instability occurs whenever $\mu\alpha < 0$
- Stoichiometry is trivial
- This is the **Keller-Segel instability (1970)**:
 - If effective interactions are attractive, aggregation instability
 - If effective interactions are repulsive, no instability
- Boring...
- **However, phenomenology explodes when we add a second species!**

Continuum theory: **binary mixtures**

- For mixtures of only two species we have

$$\mu_1 \alpha_1 \rho_{01} + \mu_2 \alpha_2 \rho_{02} < 0 \quad \text{and} \quad \delta \rho_2 = \frac{\mu_2 \rho_{02}}{\mu_1 \rho_{01}} \delta \rho_1$$

- **aggregation** if the two mobilities have **same sign**
- **separation** if the two mobilities have **opposite sign**

- Two qualitatively different cases depending on the sign of the activities:
 1. If $(\alpha_1, \alpha_2) = (+, -)$, i.e. **producer + consumer**
then instability when $\mu_2 |\alpha_2| \rho_{02} > \mu_1 |\alpha_1| \rho_{01}$
 2. If $(\alpha_1, \alpha_2) = (+, +)$, i.e. **two producers**
then instability when $\mu_2 |\alpha_2| \rho_{02} > -\mu_1 |\alpha_1| \rho_{01}$
[the case of two consumers $(\alpha_1, \alpha_2) = (-, -)$ is equivalent under the transformation $\mu_i \rightarrow -\mu_i$]

Simulations of chemically-interacting mixtures

- Far-field approximation \rightarrow superimpose the solute concentration fields generated by each colloid \rightarrow pair-wise additive velocities
- Overdamped limit \rightarrow **Brownian dynamics simulation**

$$\frac{d\mathbf{r}_i}{dt} = V_0 \sum_{k \neq i} \tilde{\alpha}_k \tilde{\mu}_i \frac{\sigma^2 \mathbf{r}_{ki}}{|\mathbf{r}_{ki}|^3} + \sum_{k \neq i} \mathbf{U}_{ik} + \boldsymbol{\xi}_i(t)$$

white noise of intensity $2D_c$

hard-sphere repulsion

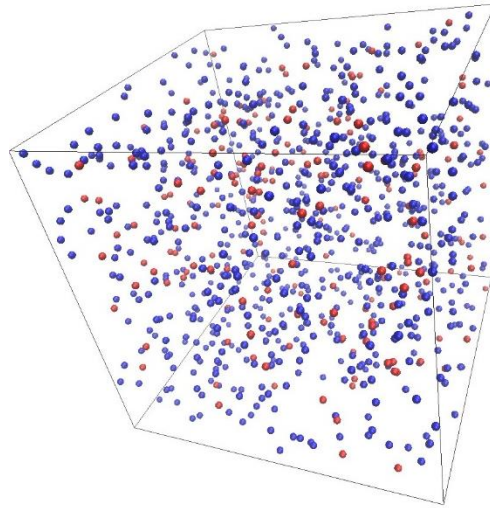
Velocity scale $V_0 = \frac{\alpha_0 \mu_0}{24\pi D}$

Noise $\tilde{D}_c = D_c / V_0 \sigma$

Activity $\tilde{\alpha} = \alpha / \alpha_0$

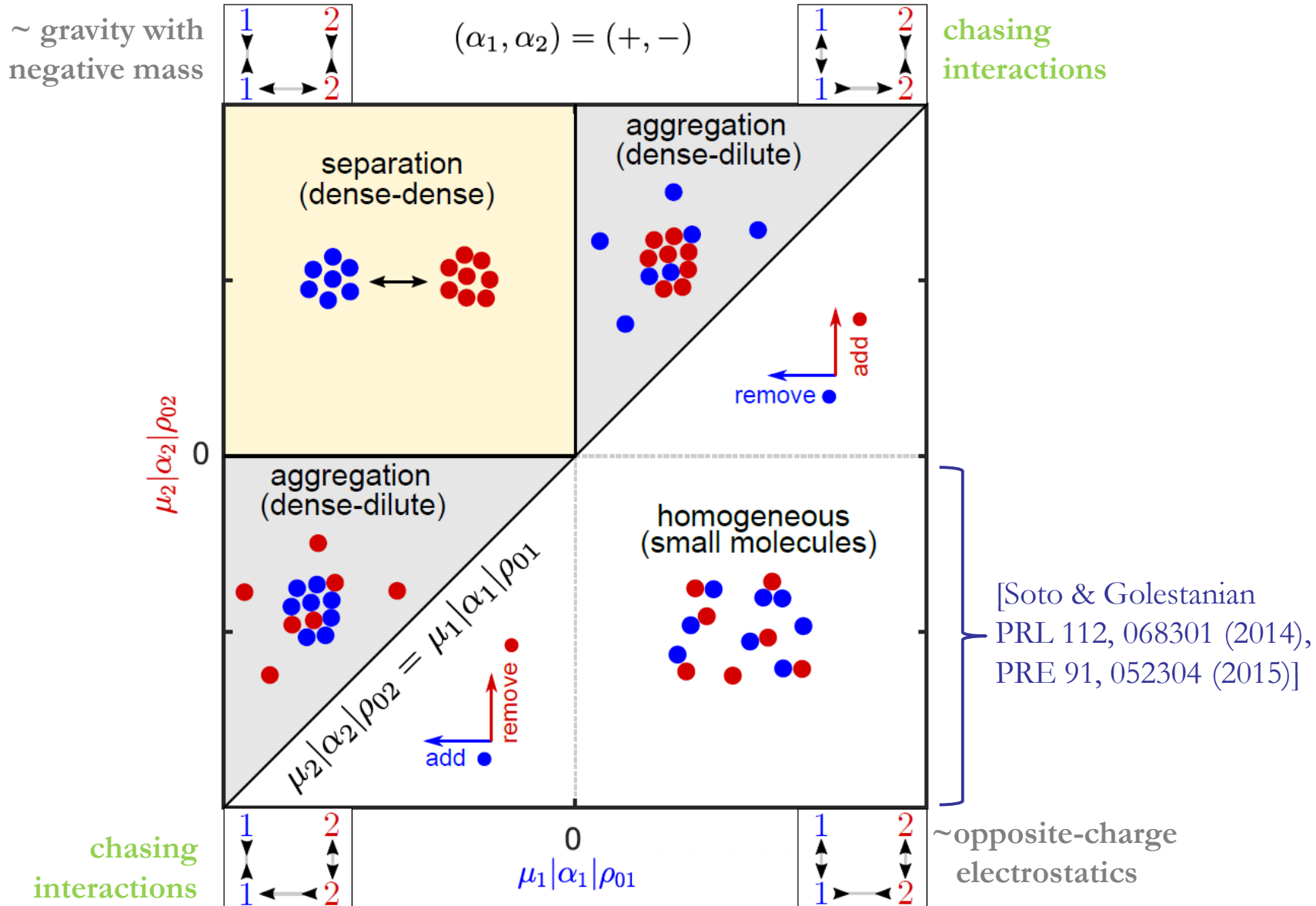
Mobility $\tilde{\mu} = \mu / \mu_0$

Simulations of chemically-interacting mixtures



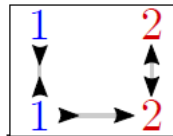
- Cubic 3D box with periodic boundary conditions:
 - Either $N_{\text{tot}} = 1000$ and $L/\sigma = 48 \rightarrow$ Volume fraction 0.0047
 - Or $N_{\text{tot}} = 4000$ and $L/\sigma = 76 \rightarrow$ Volume fraction 0.0048
- Noise $\tilde{D}_c = 0.01$
- Time step $\delta t = 0.001\sigma/V_0$
- Run time between $5 \cdot 10^4$ and $8 \cdot 10^5$ steps

Phase diagram for mixtures of **producer + consumer**

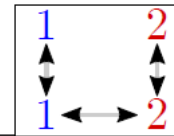


Phase diagram for mixtures of two producers

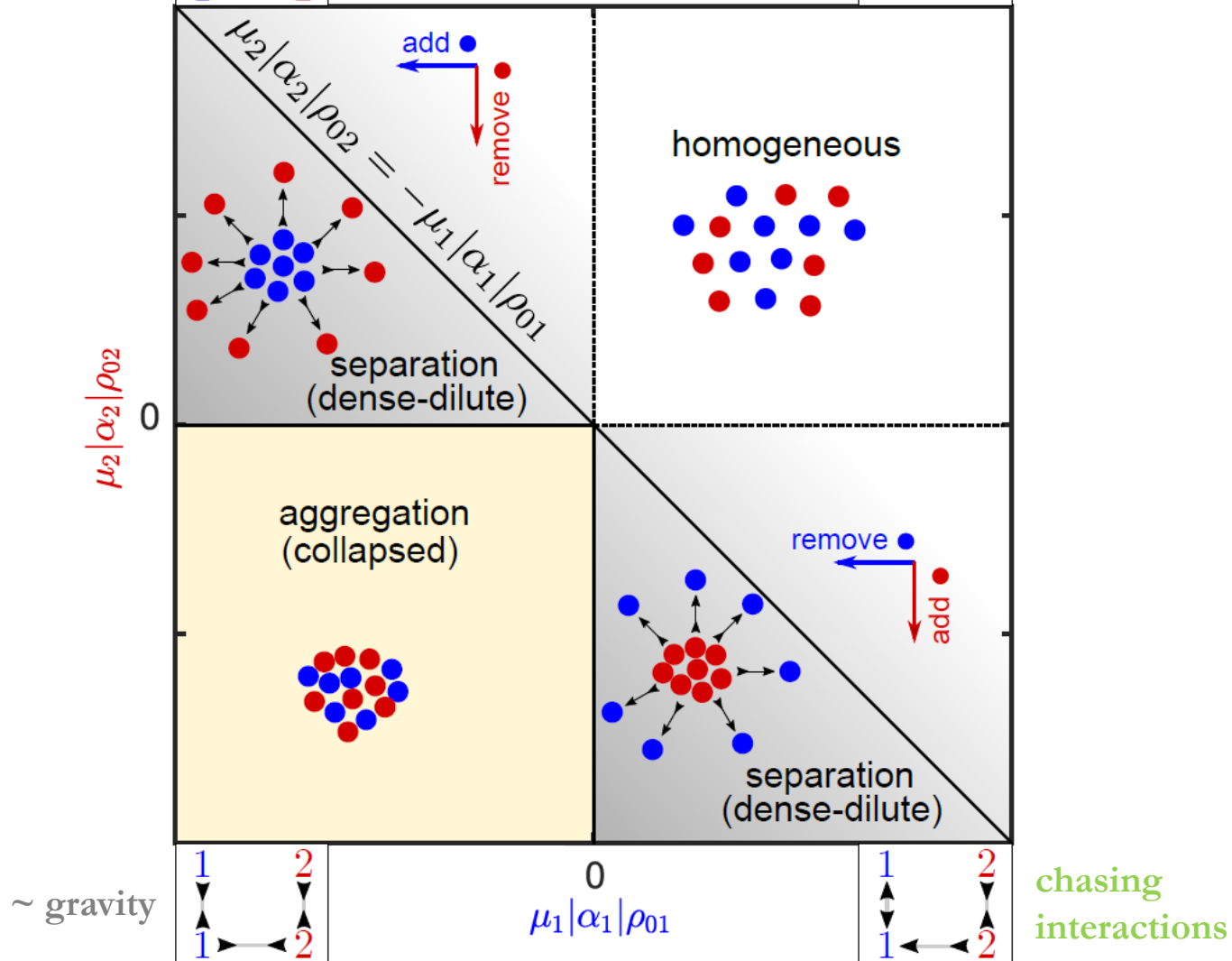
chasing interactions



$$(\alpha_1, \alpha_2) = (+, +)$$



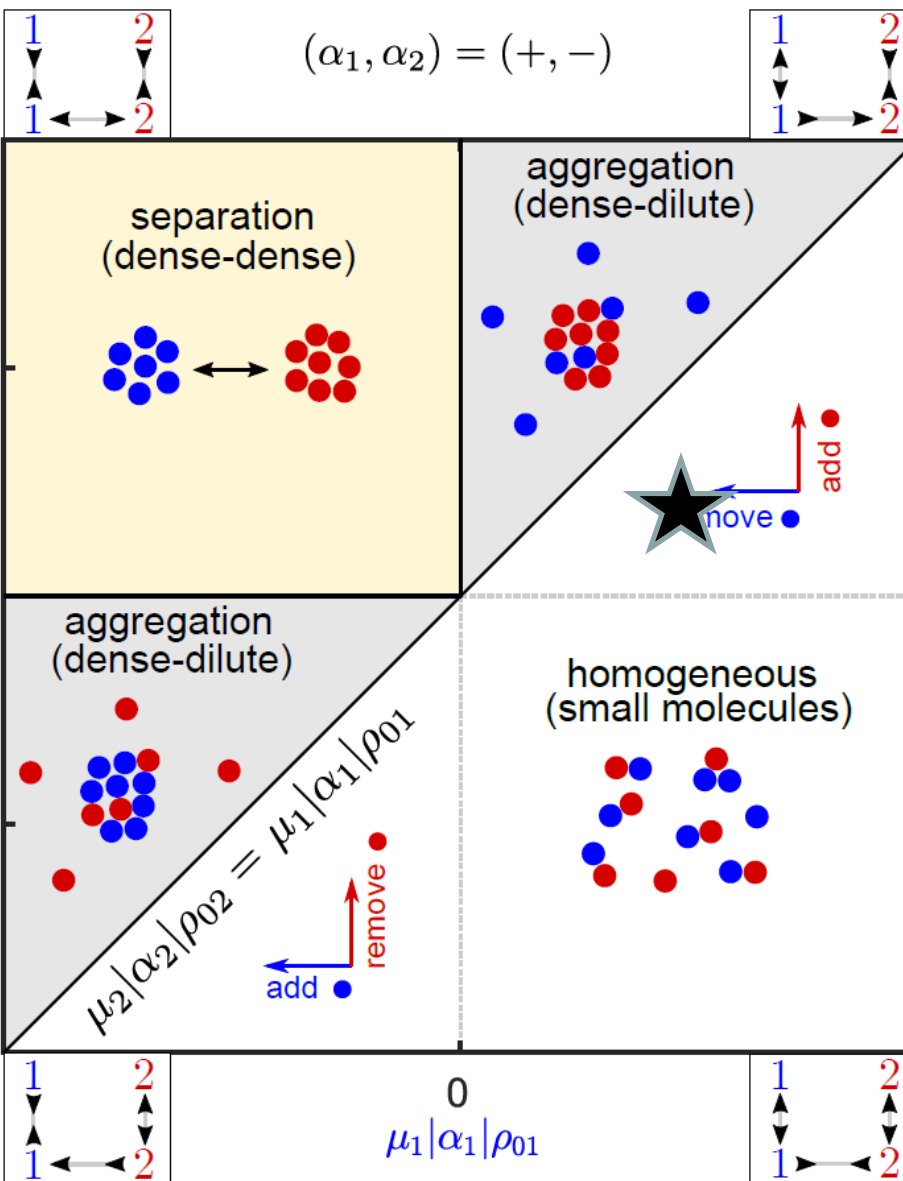
~like-charge electrostatics



Producer+consumer:

homogeneous state with active molecule formation

~ gravity with negative mass



Producer+consumer:

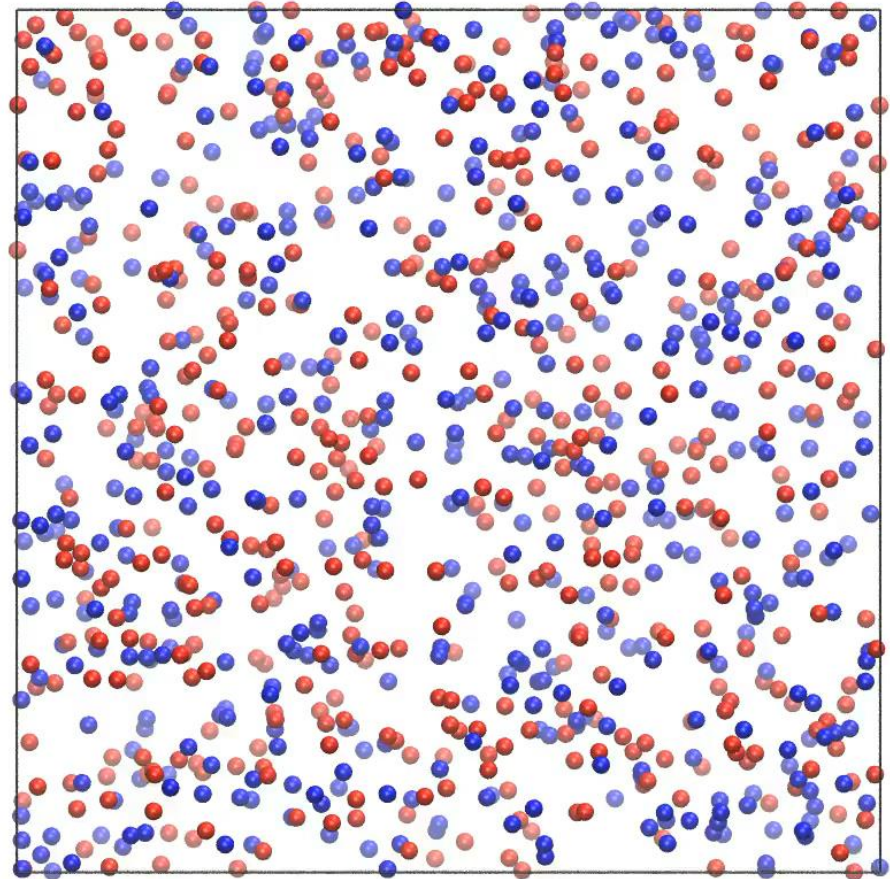
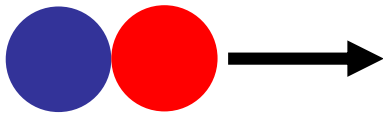
homogeneous state with active molecule formation

$$\tilde{\alpha} = (1, -1)$$

$$\tilde{\mu} = (1, 1/2)$$

$$N = (500, 500)$$

In this case, formation of
self-propelled dimers



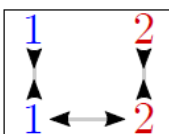
3D realisation of the 2D active molecules explored in

[Soto, R., & Golestanian, R. Phys. Rev. Lett 112, 068301 (2014), Phys. Rev. E 91, 052304 (2015)]

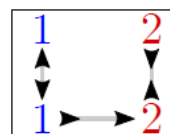
Producer+consumer:

aggregation into cluster with defined stoichiometry

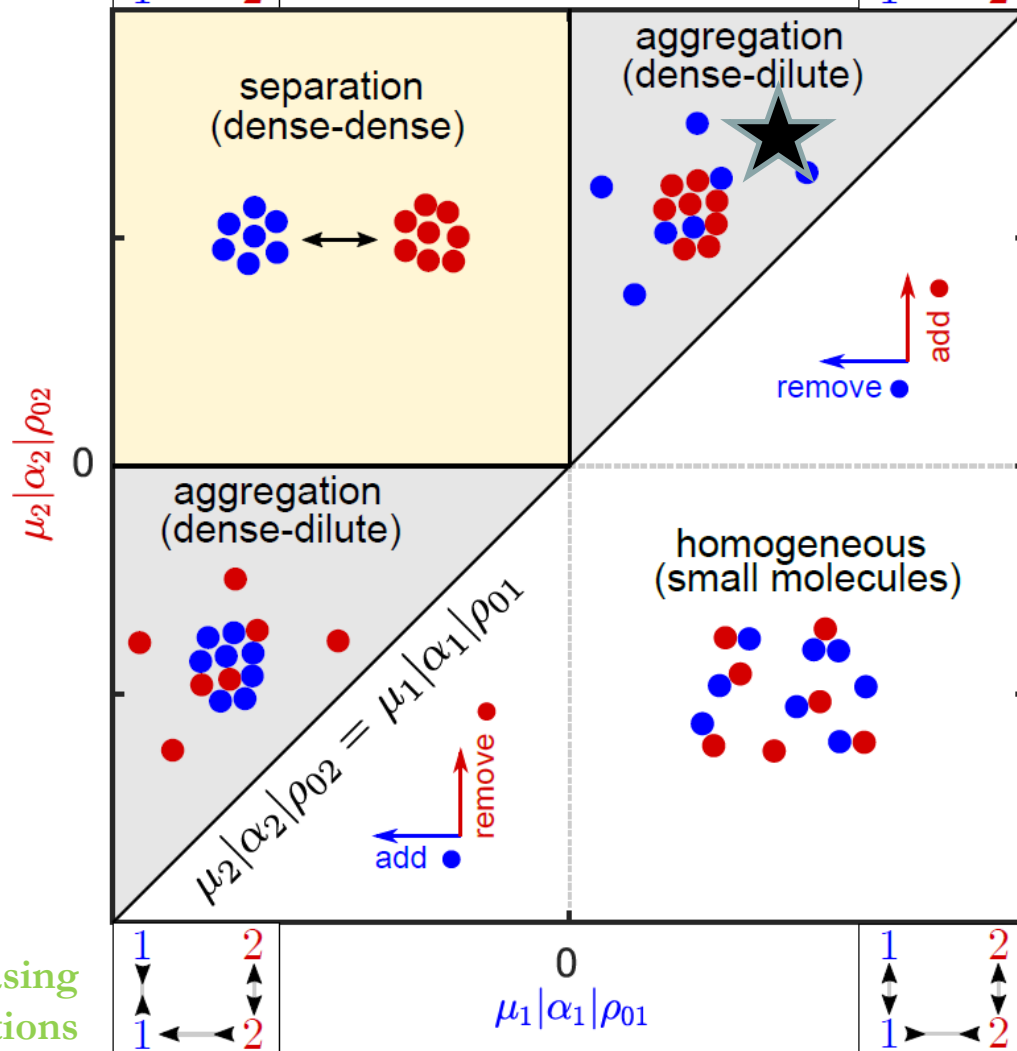
~ gravity with
negative mass



$$(\alpha_1, \alpha_2) = (+, -)$$



chasing
interactions



[Soto & Golestanian
PRL 112, 068301 (2014),
PRE 91, 052304 (2015)]

chasing
interactions

~ opposite-charge
electrostatics

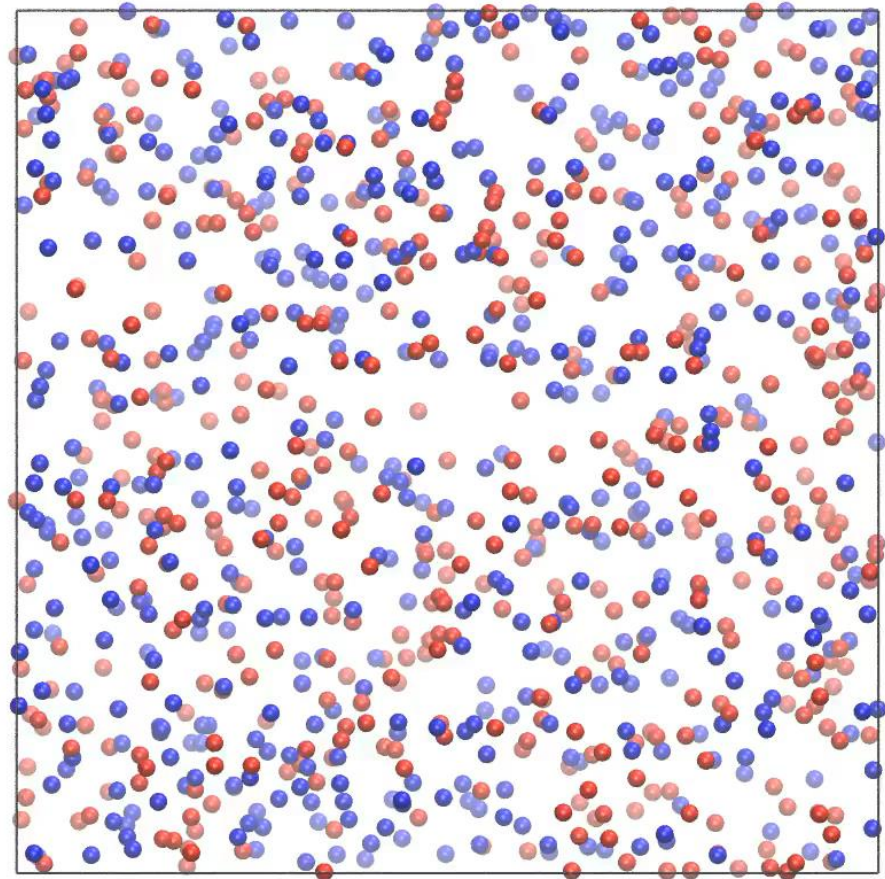
Producer+consumer:

aggregation into cluster with defined stoichiometry

$$\tilde{\alpha} = (1, -1/2)$$

$$\tilde{\mu} = (1, 6)$$

$$N = (500, 500)$$



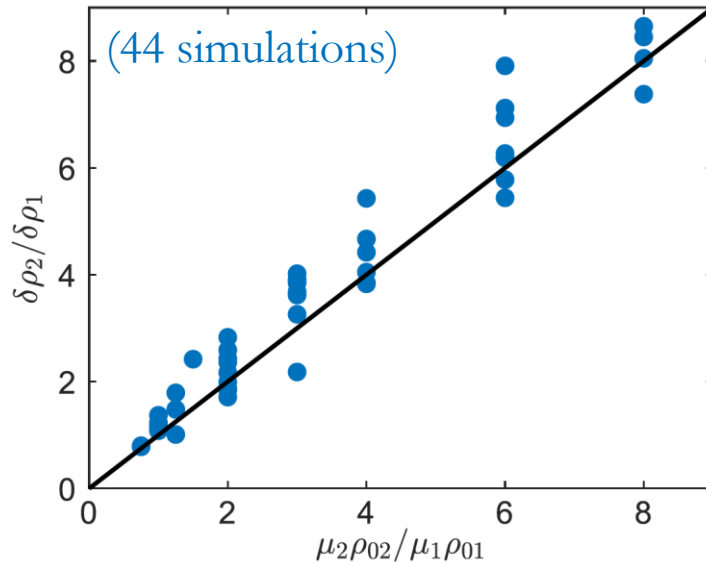
Producer+consumer:

aggregation into cluster with defined stoichiometry

- The initial stoichiometry of the simulations is well captured by the continuum theory prediction

$$\delta\rho_2 = \frac{\mu_2\rho_{02}}{\mu_1\rho_{01}}\delta\rho_1$$

(for a binary mixture)



**Initial stoichiometry
depends only on mobility,
independent of activity**

Producer+consumer:

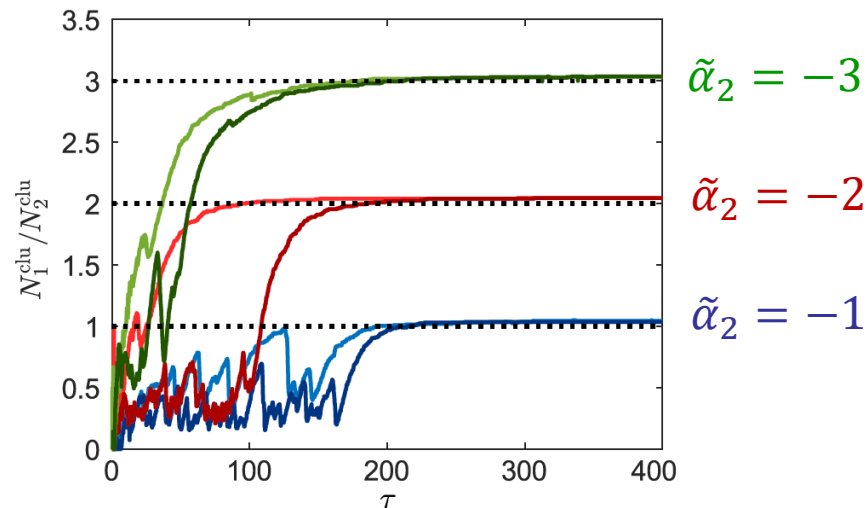
aggregation into cluster with defined stoichiometry

- The final stoichiometry of the cluster is determined by:
 - All particles of the self-attractive type are in the cluster

$$N_2^{\text{clu}} = N_2$$

- Particles of the self-repelling type are added until neutrality is reached

$$\alpha_1 N_1^{\text{clu}} + \alpha_2 N_2^{\text{clu}} = 0$$



**Final stoichiometry
depends only on activity,
independent of mobility**

- **Metabolon formation: all product of one enzyme is channeled in the exact amount as substrate to the next enzyme**

Producer+consumer:

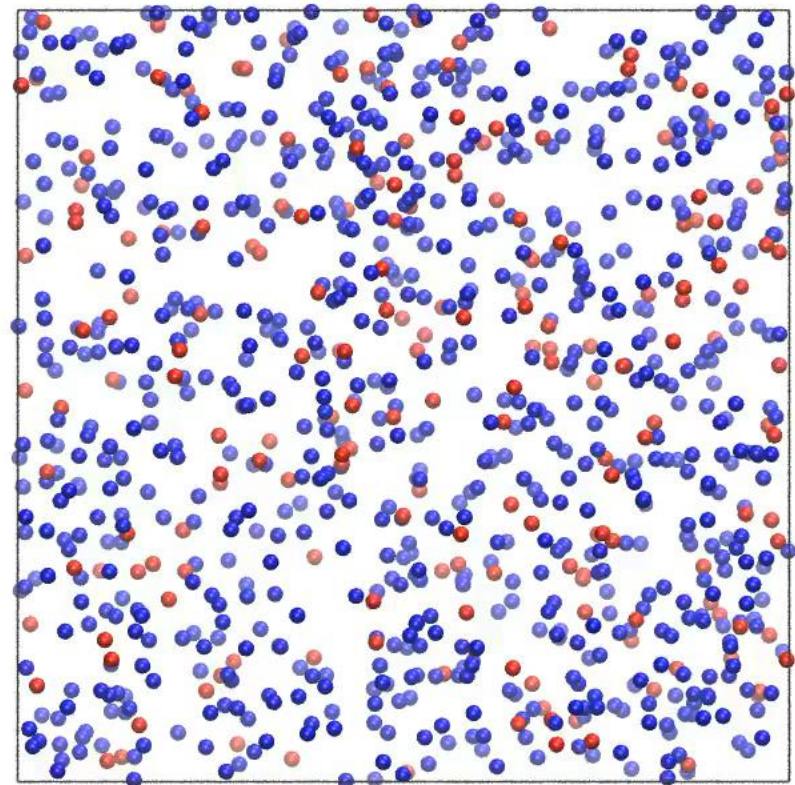
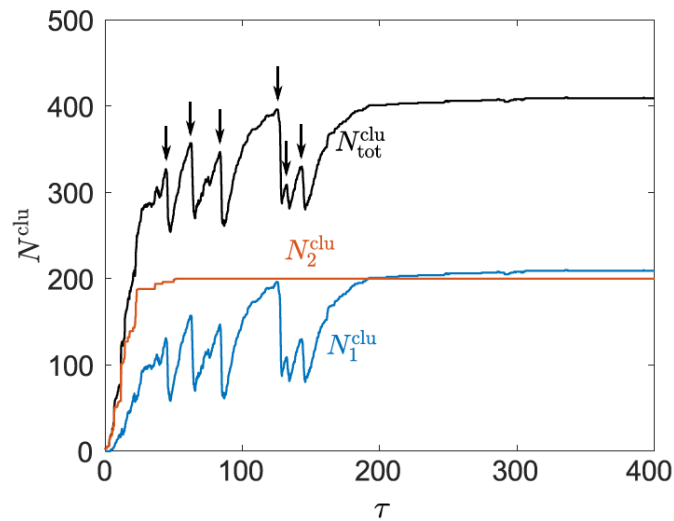
aggregation into cluster with defined stoichiometry
→ shape-instability towards self-propelled state

$$\tilde{\alpha} = (1, -1)$$

$$\tilde{\mu} = (1, 8)$$

$$N = (800, 200)$$

Bistability between static and self-propelled



Self-propelled states typically observed when $|\alpha_2|\mu_2 \gg |\alpha_1|\mu_1$

Producer+consumer:

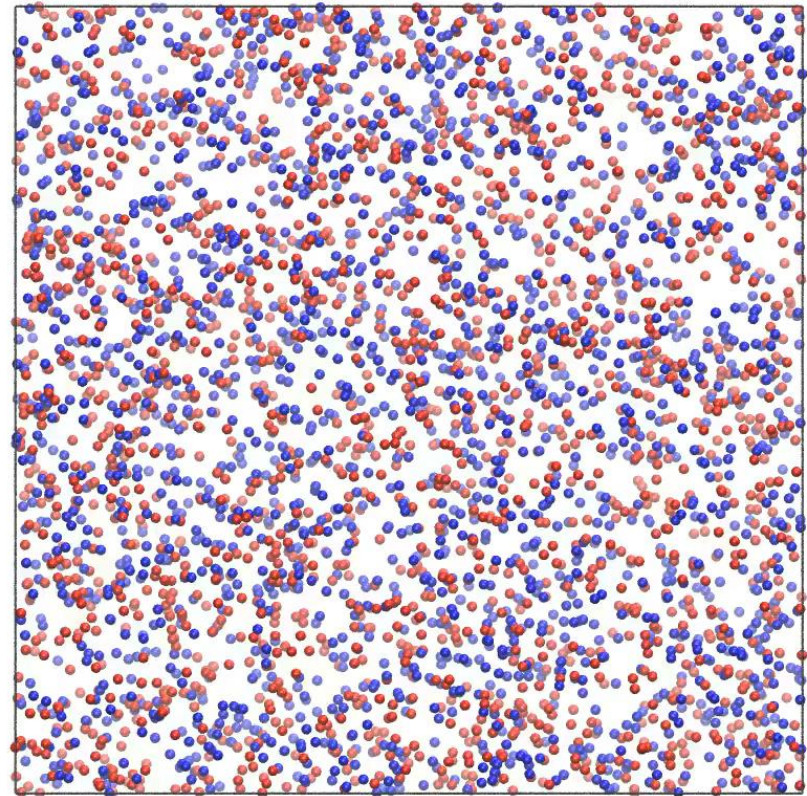
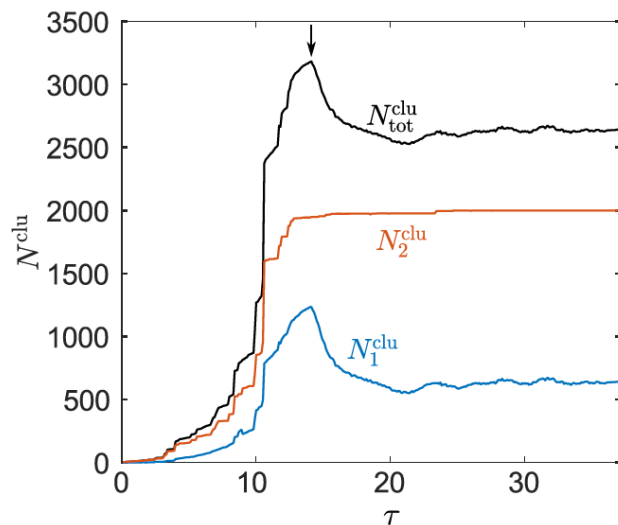
aggregation into cluster with defined stoichiometry
→ shape-instability towards self-propelled state

$$\tilde{\alpha} = (1, -2)$$

$$\tilde{\mu} = (1, 3)$$

$$N = (2000, 2000)$$

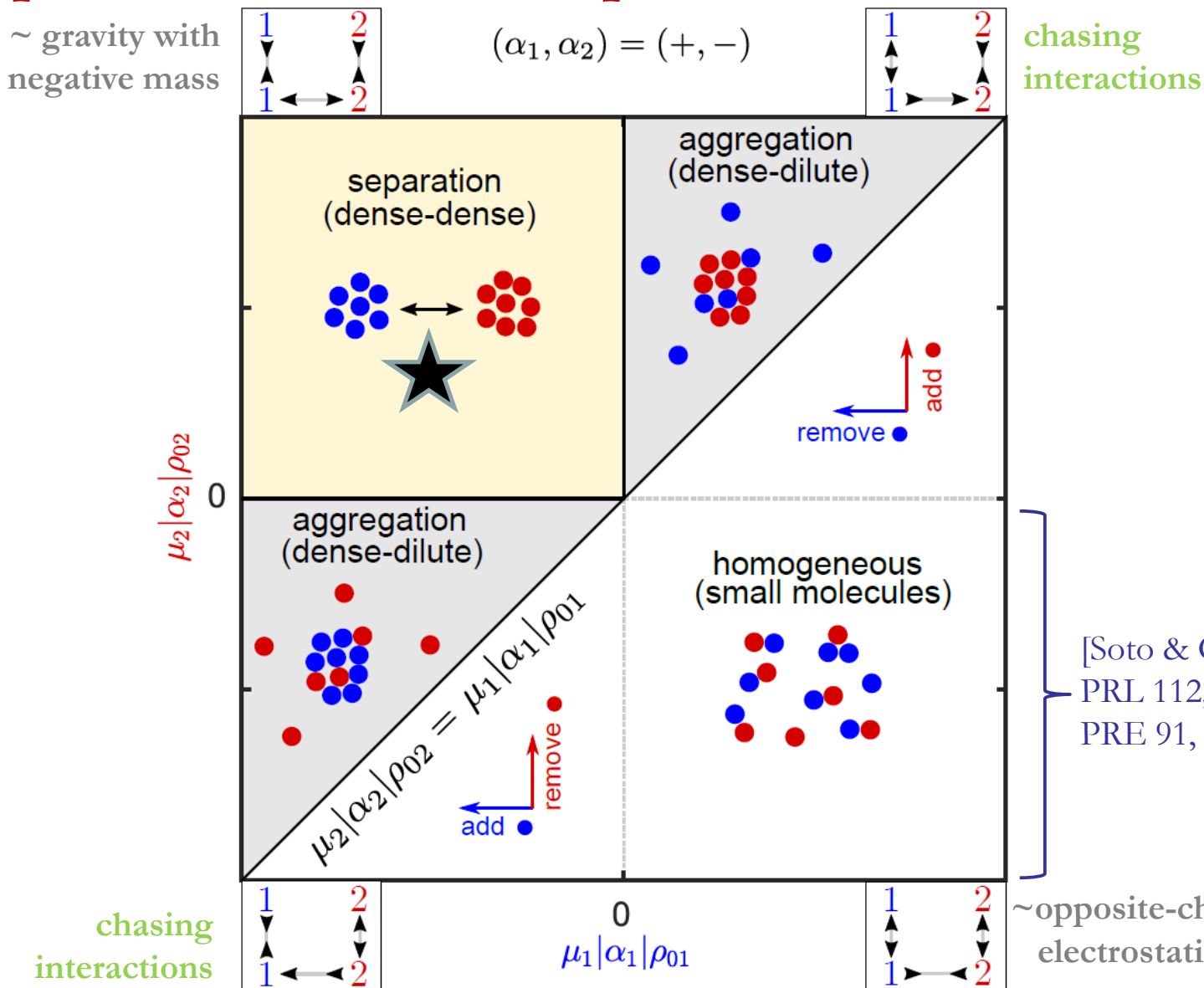
Stable self-propelled state



► Morphogenesis and collective migration: small changes in a subpopulation can trigger collective self-propulsion

Producer+consumer:

separation into two collapsed clusters



[Soto & Golestanian
PRL 112, 068301 (2014),
PRE 91, 052304 (2015)]

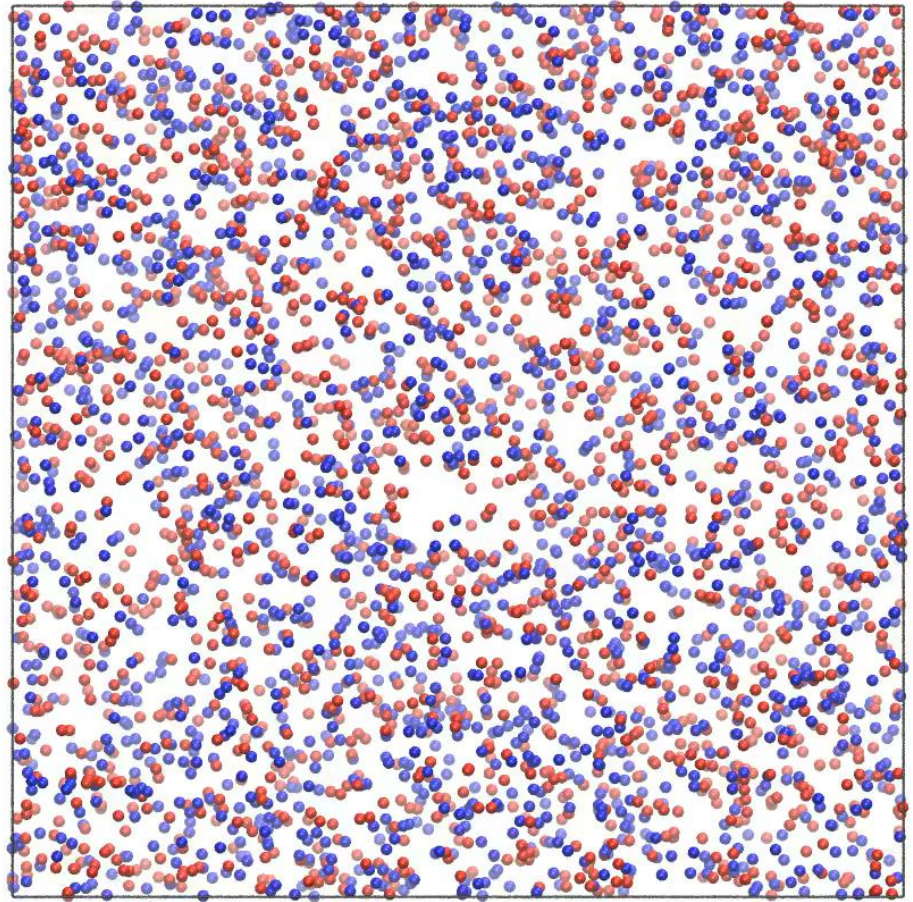
Producer+consumer:

separation into two collapsed clusters

$$\tilde{\alpha} = (1, -1)$$

$$\tilde{\mu} = (-1, 2)$$

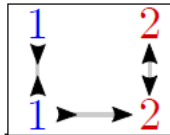
$$N = (2000, 2000)$$



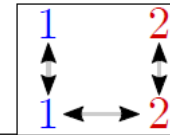
Two producers:

homogeneous state without molecule formation

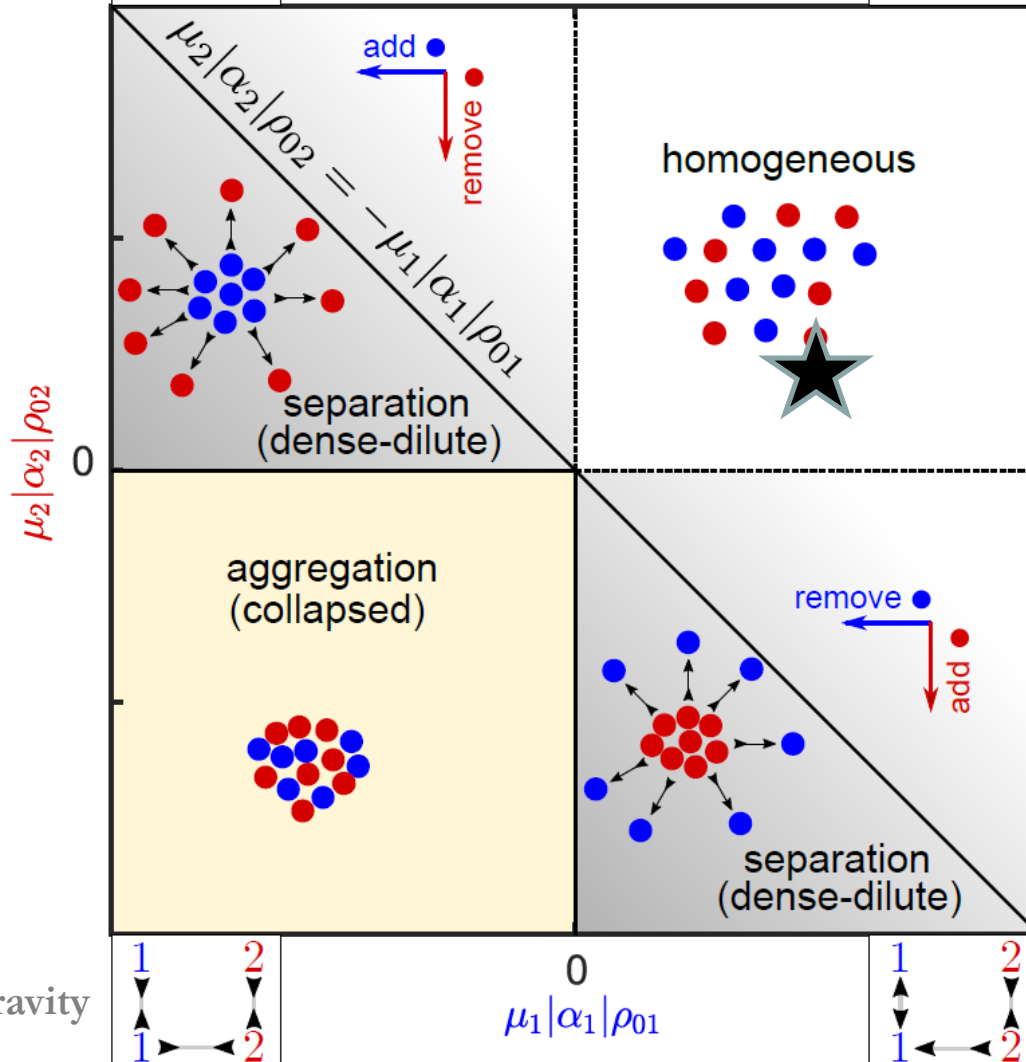
chasing interactions



$$(\alpha_1, \alpha_2) = (+, +)$$



~like-charge electrostatics



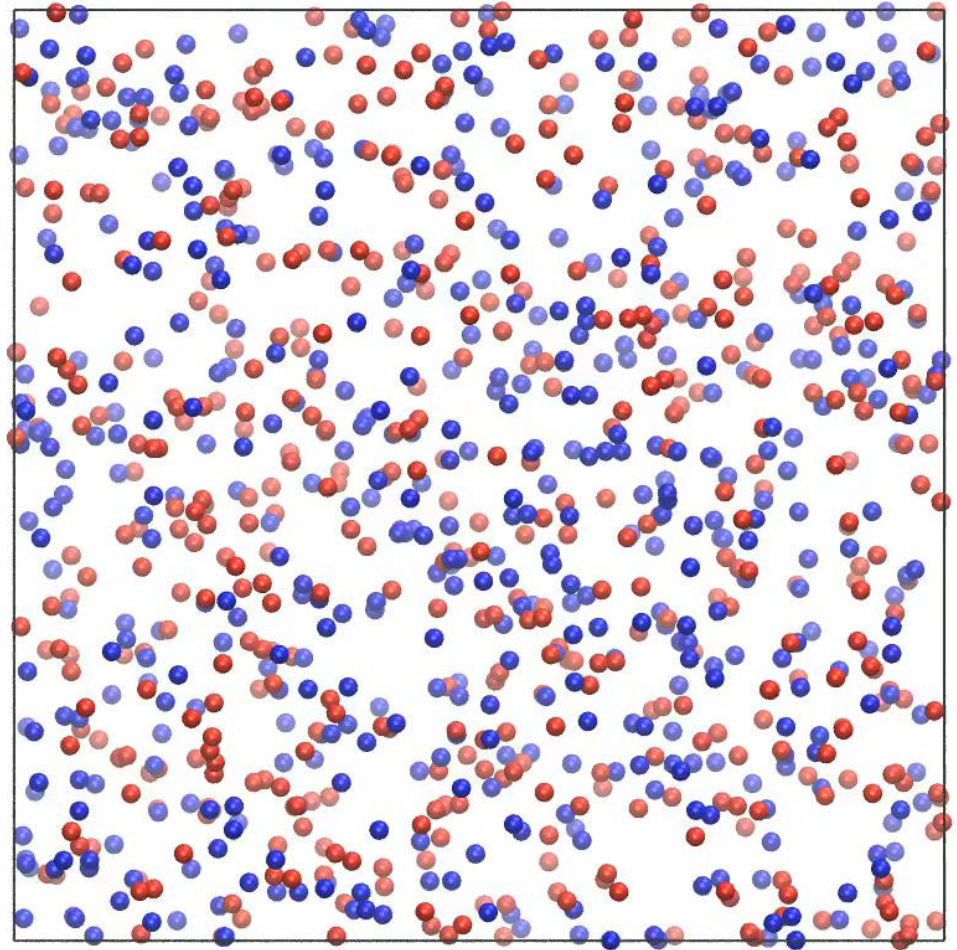
Two producers:

homogeneous state without molecule formation

$$\tilde{\alpha} = (1,1)$$

$$\tilde{\mu} = (1,1/2)$$

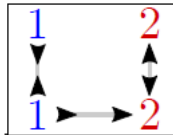
$$N = (500,500)$$



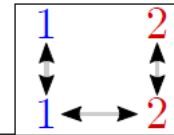
Two producers:

aggregation into a collapsed cluster

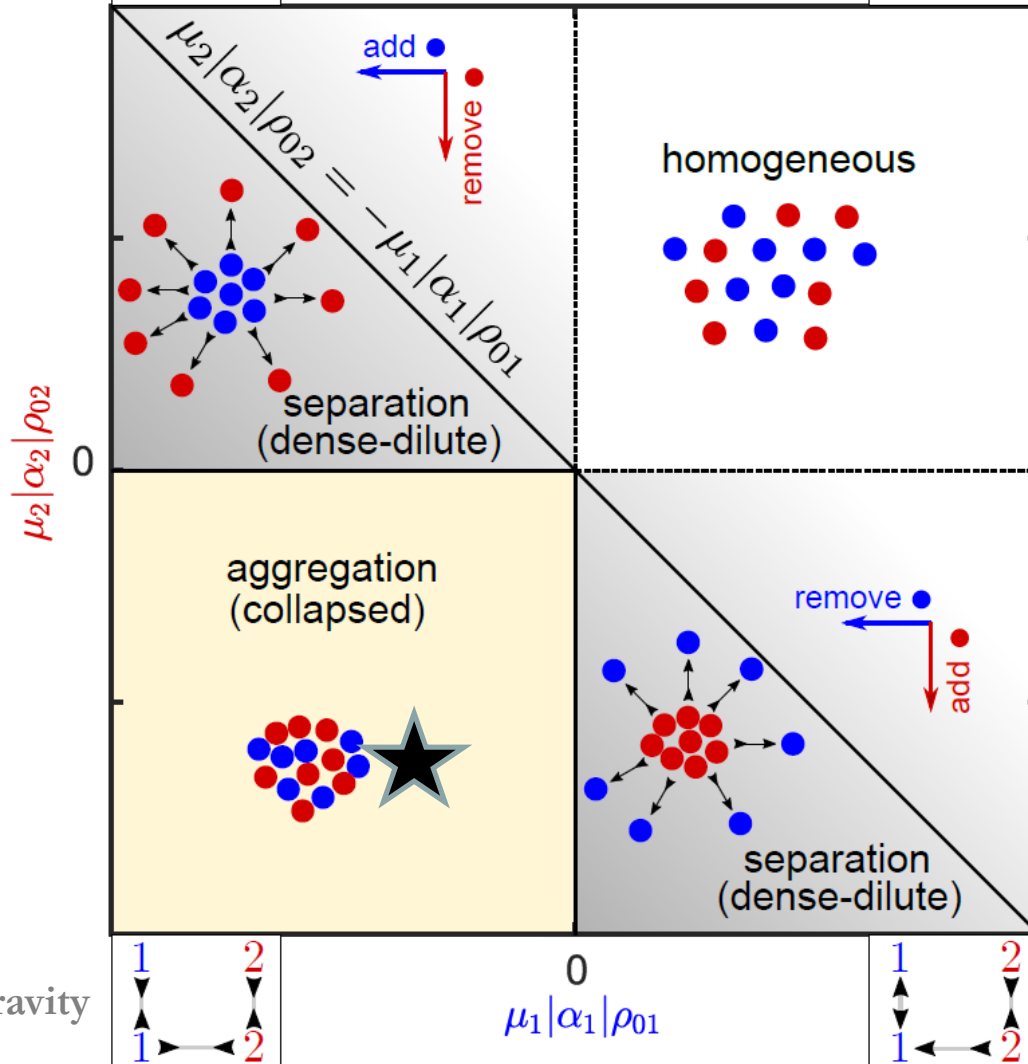
chasing interactions



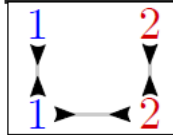
$$(\alpha_1, \alpha_2) = (+, +)$$



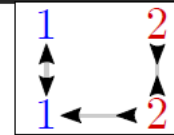
~like-charge electrostatics



~ gravity



$$\mu_1 |\alpha_1| \rho_1$$



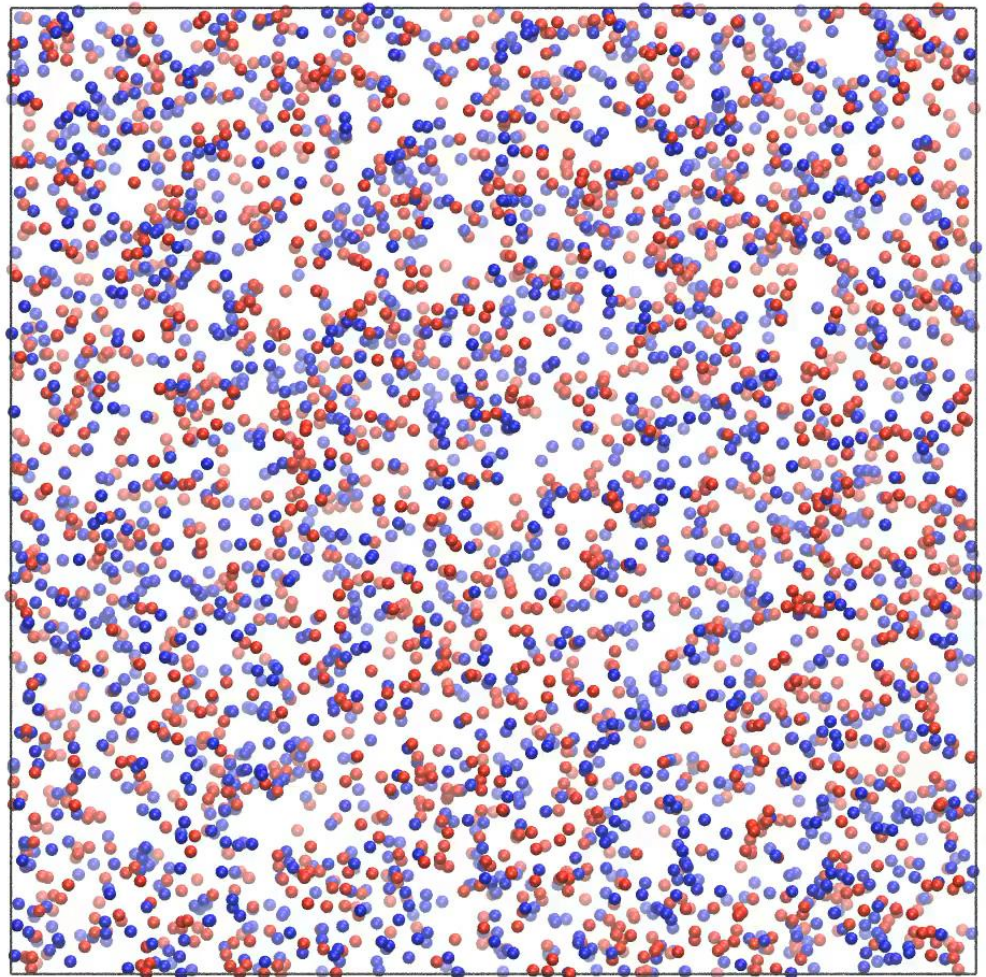
chasing interactions

Two producers:
aggregation into a collapsed cluster

$$\tilde{\alpha} = (1,1)$$

$$\tilde{\mu} = (-1,-2)$$

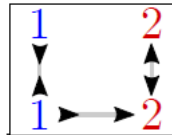
$$N = (2000,2000)$$



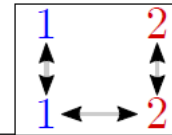
Two producers:

separation into cluster and repelled dilute phase

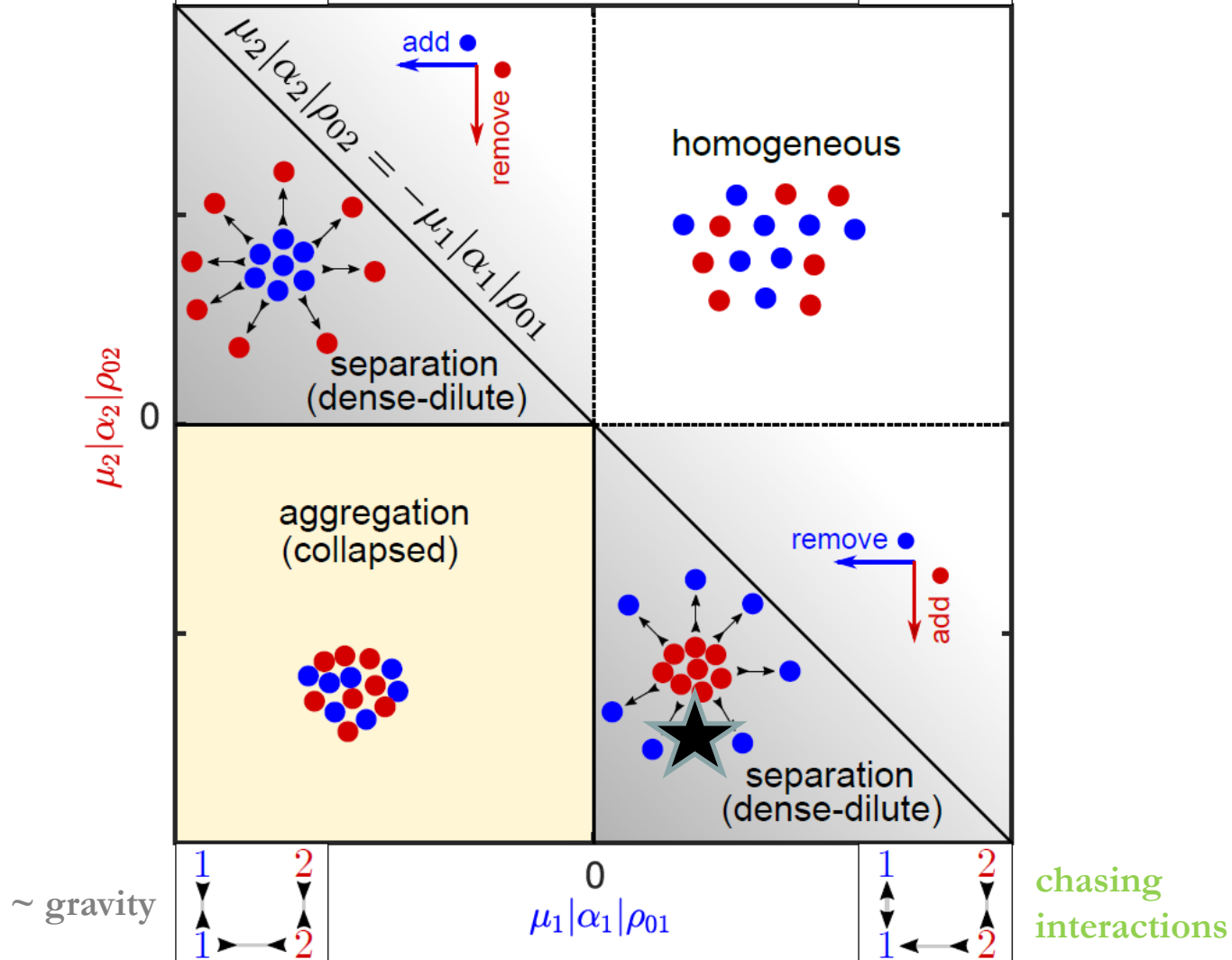
chasing interactions



$$(\alpha_1, \alpha_2) = (+, +)$$



~like-charge electrostatics



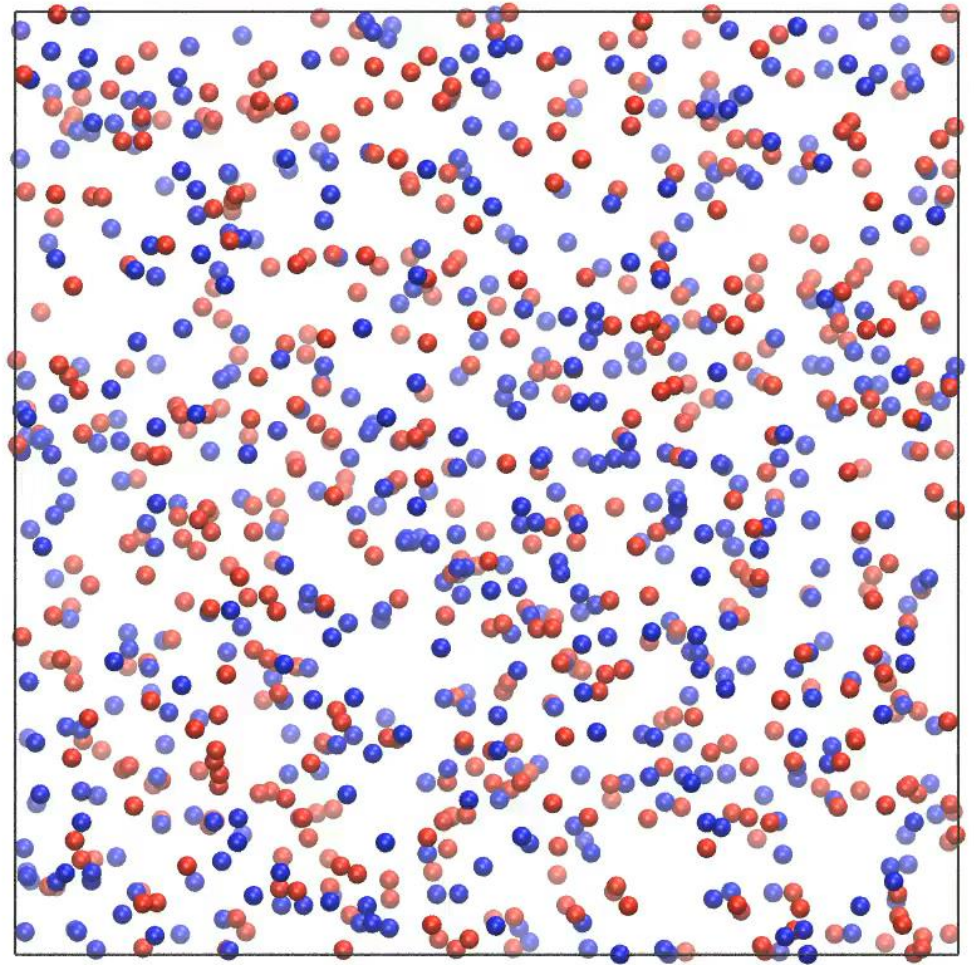
Two producers:

separation into cluster and repelled dilute phase

$$\tilde{\alpha} = (1,1)$$

$$\tilde{\mu} = (1,-2)$$

$$N = (500,500)$$



Beyond binary mixtures:

phase separation triggered by active “doping agent”

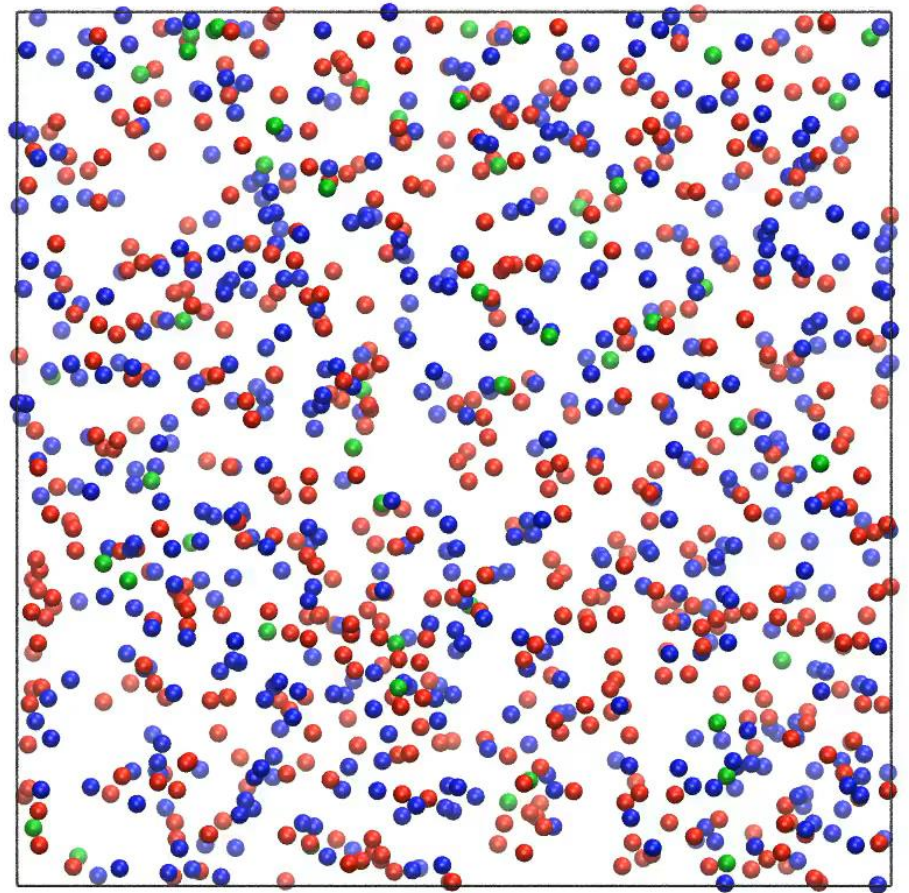
$$\tilde{\alpha} = (1, -1, -5)$$

$$\tilde{\mu} = (1, 1/2, 2)$$

$$N = (500, 500, 50)$$

Addition of just a few particles of a 3rd type triggers phase separation in an otherwise homogeneous mixture (cf. first movie)

$$\sum_i \mu_i \alpha_i \rho_{0i} < 0$$

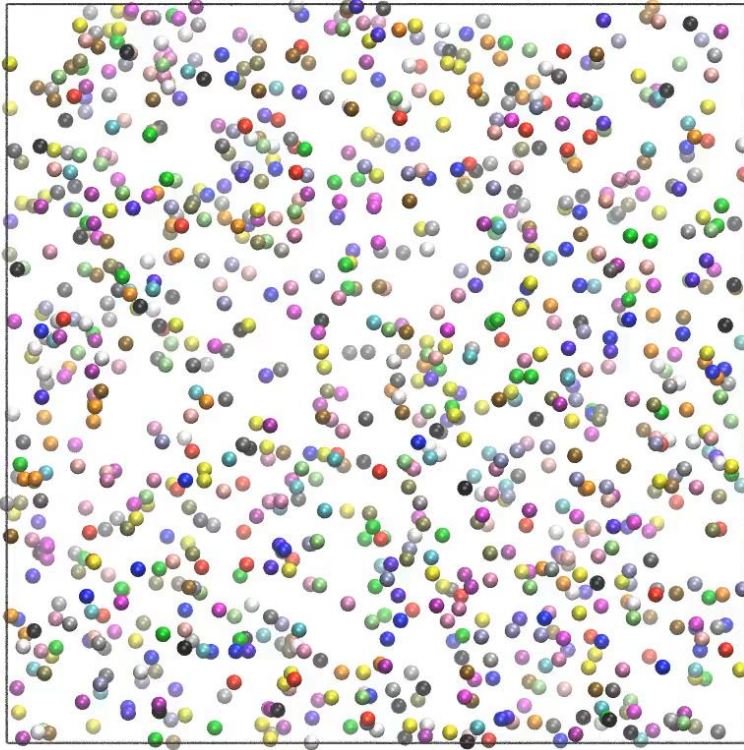


Beyond binary mixtures:

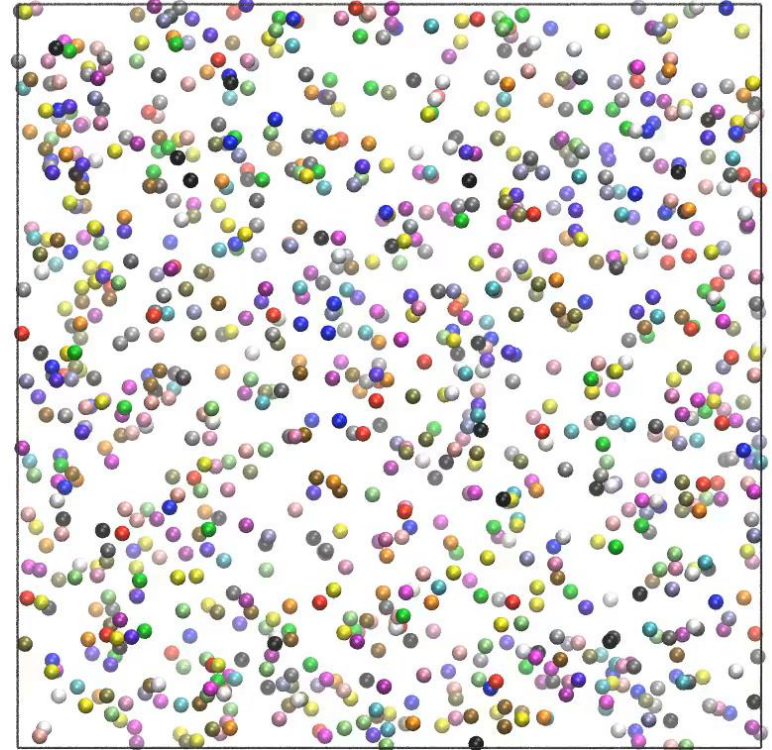
phase separation in highly polydisperse mixtures

- Randomly-generated mixtures with 20 different species

homogeneous



phase-separating



Beyond binary mixtures:

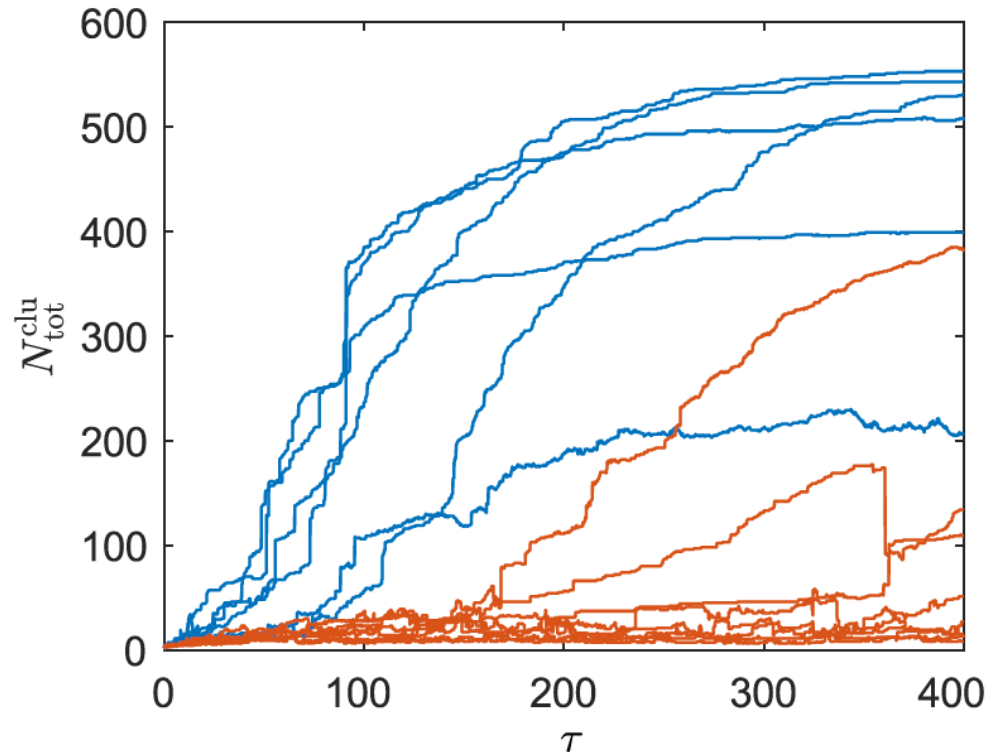
phase separation in highly polydisperse mixtures

- Randomly-generated mixtures with 20 different species

Instability condition:

$$\sum_i \mu_i \alpha_i \rho_{0i} < 0$$

All **linearly unstable** mixtures undergo phase separation, but some **linearly stable** mixtures too
→ nucleation-and-growth mechanism?



Summary II:

- Theory for **arbitrary mixtures** (N species) of chemically-interacting particles shows **wide variety** of phase separation phenomena
- Linked to **active, nonreciprocal** nature of the interactions
- Minimal model applicable to catalytic **enzymes** (metabolon formation), heterogeneous populations of **cells and bacteria** (morphogenesis, migration), catalytic phoretic **colloids** (self-assembly of active materials)

[Agudo-Canalejo, J., & Golestanian, R. (2019). Active phase separation in mixtures of chemically interacting particles. *Physical Review Letters*, 123, 018101.]

